

Relativistic Shifts of g_μ in Muonic Atoms

Jess H. Brewer

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The discovery of an M9B tune giving substantial transverse muon spin polarization allows convenient μ^- SR experiments in high transverse magnetic field. This new capability has made possible precise measurements of the magnetogyric ratios of negative muons in the ground states of muonic atoms of ^{12}C , ^{16}O , ^{24}Mg , ^{28}Si , ^{32}S , ^{40}Ca , $^{\text{nat}}\text{Ti}$, $^{\text{nat}}\text{Zn}$, $^{\text{nat}}\text{Cd}$ and $^{\text{nat}}\text{Pb}$. The precision for $^{12}\text{C}\mu^-$ is ± 23 ppm, of which only 6 ppm is statistical; for $^{\text{nat}}\text{Zn}\mu^-$ the precision is ± 269 ppm and for $^{\text{nat}}\text{Pb}\mu^-$ it is $\pm 0.23\%$. Such experiments may eventually provide a new testing ground for QED in very strong Coulomb fields; today they offer a new way of measuring finite nuclear size effects. I will discuss a little of the history of TRIUMF Experiment 932 and explain how these measurements emerged “accidentally” from an attempt to improve the performance of μ^- SR experiments in general.

Phys. Rev. A Brief Report, in press (2005)

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Facility & Method used:

μ SR *rotation*
relaxation
resonance

<i>m</i>	<i>s</i>	Applied*
<i>u</i>	<i>p</i>	Elementary
<i>o</i>	<i>i</i>	Particle
<i>n</i>	<i>n</i>	Physics

*(to basic research in
Materials Science
and Chemistry)

[and “Fundamental” Physics]

Visit <http://musr.org>

"Themes" in $\mu^\pm SR$

μ^+ only (?)

Muonium as light Hydrogen

($Mu = \mu^+ e^-$)

($H = p^+ e^-$)

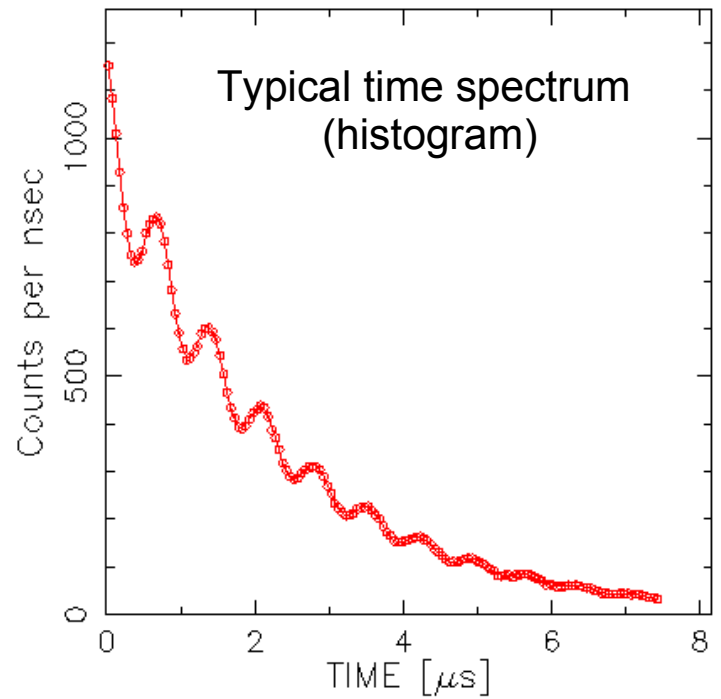
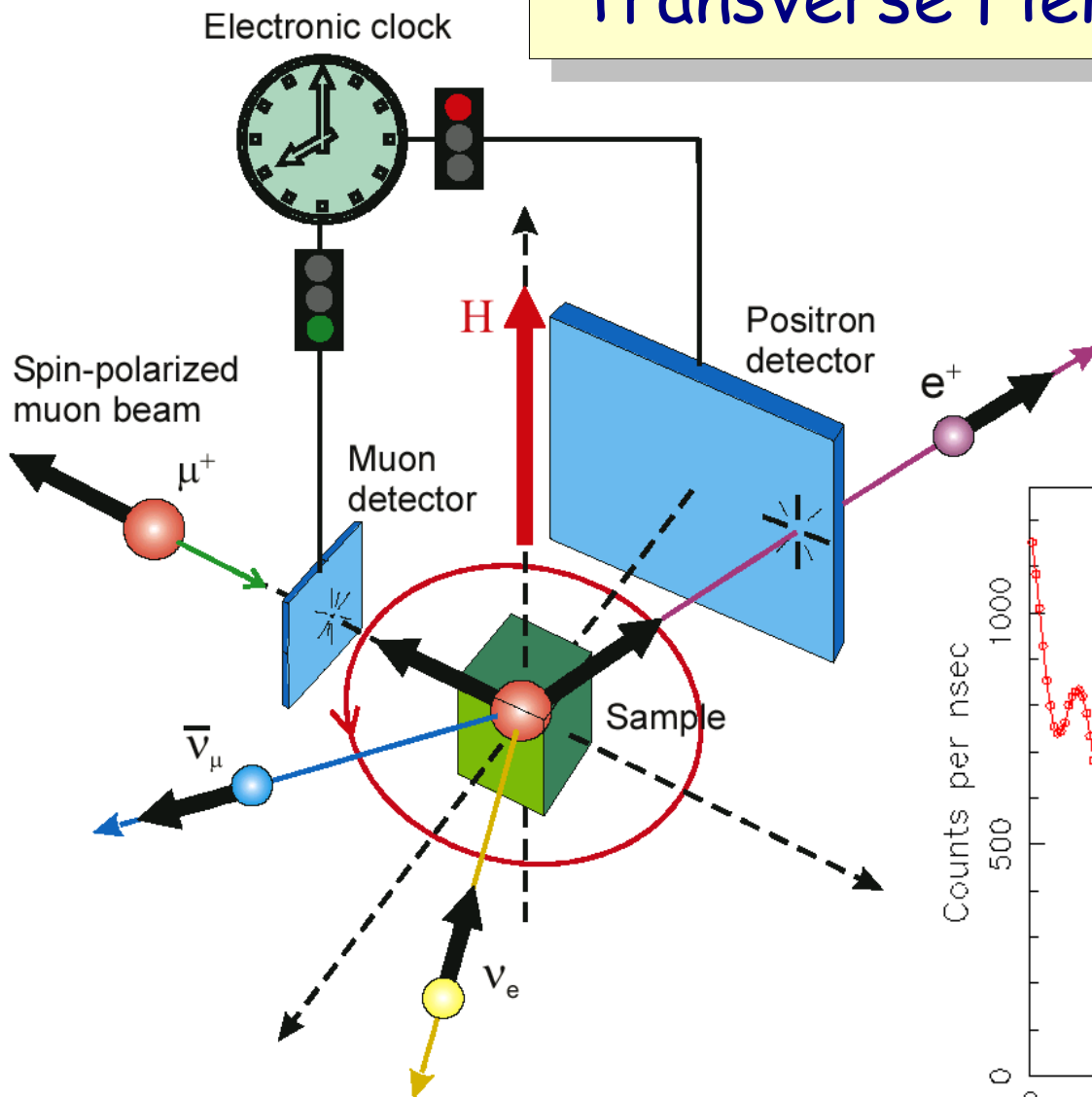
- **Mu vs. H atom Chemistry:**
 - *gases, liquids & solids*
 - *Best test of reaction rate theories.*
 - *Study "unobservable" H atom rxns.*
 - *Discover new radical species.*
- **Mu vs. H in Semiconductors:**
 - *Until recently, $\mu^+ SR \rightarrow$ only data on metastable H states in semiconductors!*
- **Quantum Diffusion:** μ^+ in metals (compare H^+); **Mu** in nonmetals (compare H).
- **Ultra-Heavy Hydrogen:** neutral muonic helium ($\alpha^{++} \mu^- e^-$) has $m \approx 4.11 m_H$
- **Muonic Atoms:** HF transitions with nucleus; relativistic shifts of g_μ

μ^+ or μ^-

The Muon as a Probe

- **Probing Magnetism:** unequalled sensitivity
 - Local fields: electronic structure; ordering
 - Dynamics: electronic, nuclear spins
- **Probing Superconductivity:** (esp. $HT_c SC$)
 - Coexistence of SC & Magnetism
 - Magnetic Penetration Depth λ
 - Coherence Length ξ

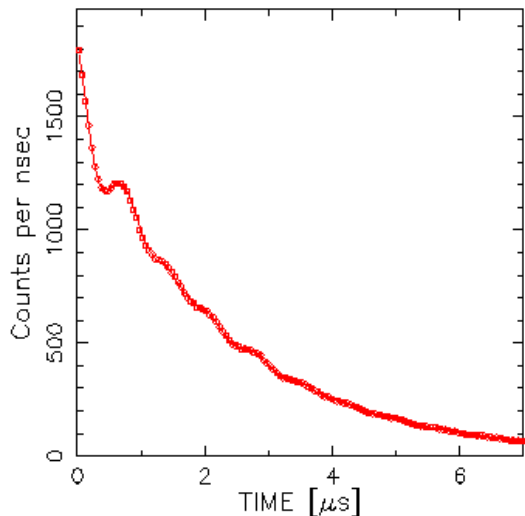
Transverse Field (TF) μ^+SR



μ^+SR

vs.

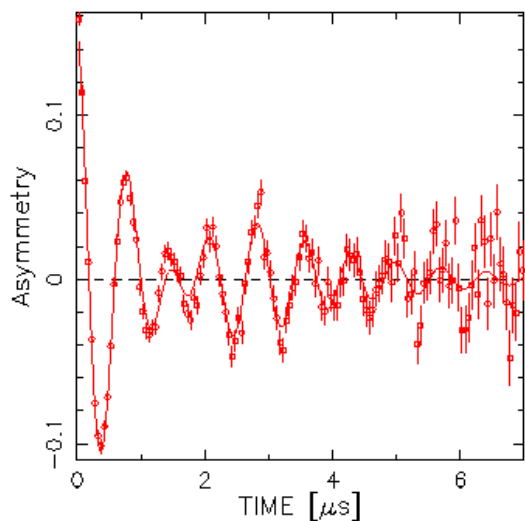
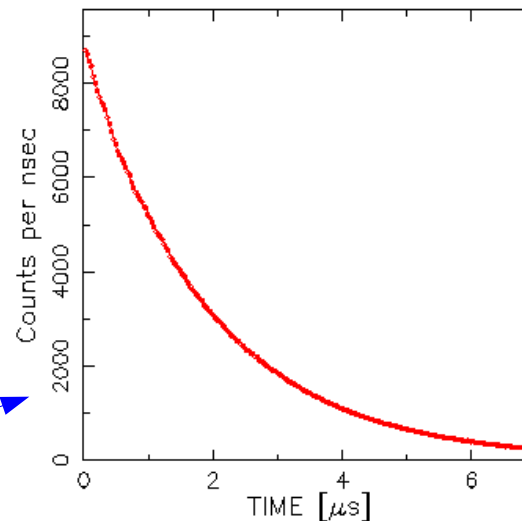
μ^-SR



Typical time spectrum (histogram)

Single lifetime
 $\tau_\mu = 2.197 \mu s$

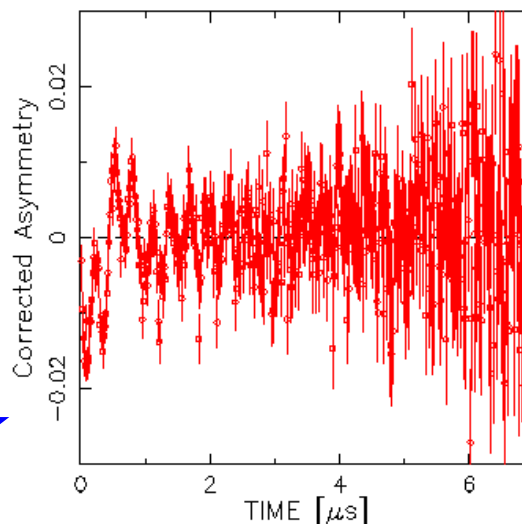
Multiple lifetimes (some very short!)



Asymmetry spectrum

Large amplitudes

Small amplitudes



TRIUMF Experiment # **932**:

Improving μ^-SR Performance





Initial proposal (Dec 2001) to M.S. EEC:
M9B at **LOW** priority(!)

We take **only** beam **no one else wants**.

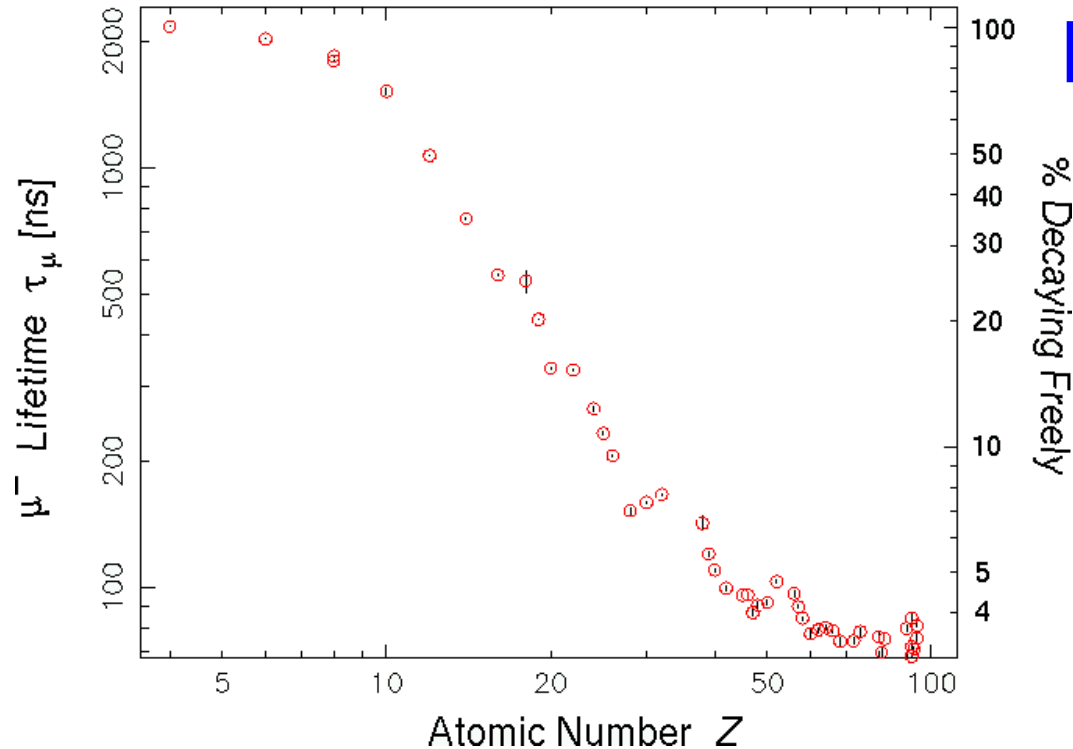
μ^-SR : It is easy to get the impression that **only positive** muons are employed in μSR .

Although *most* μSR is μ^+SR , it is often desirable to use **negative** muons in the same way, albeit with more **difficulty**. It is the goal of **E932** to **reduce** that **difficulty**.

Drawbacks of μ^-SR → Proposed Mitigations

- Nuclear Muon Capture: short lifetimes, few decay e^- → • Look for neutron asymmetries in heavier elements 
- **$L \cdot S$** Depolarization in the atomic cascade → • "Tag" events with specific muonic X-rays 
- Giant Hyperfine Int. with nonzero-spin nuclei → • Observe characteristic F^\pm precession signals 

Negative Muon Lifetimes in Muonic Atoms



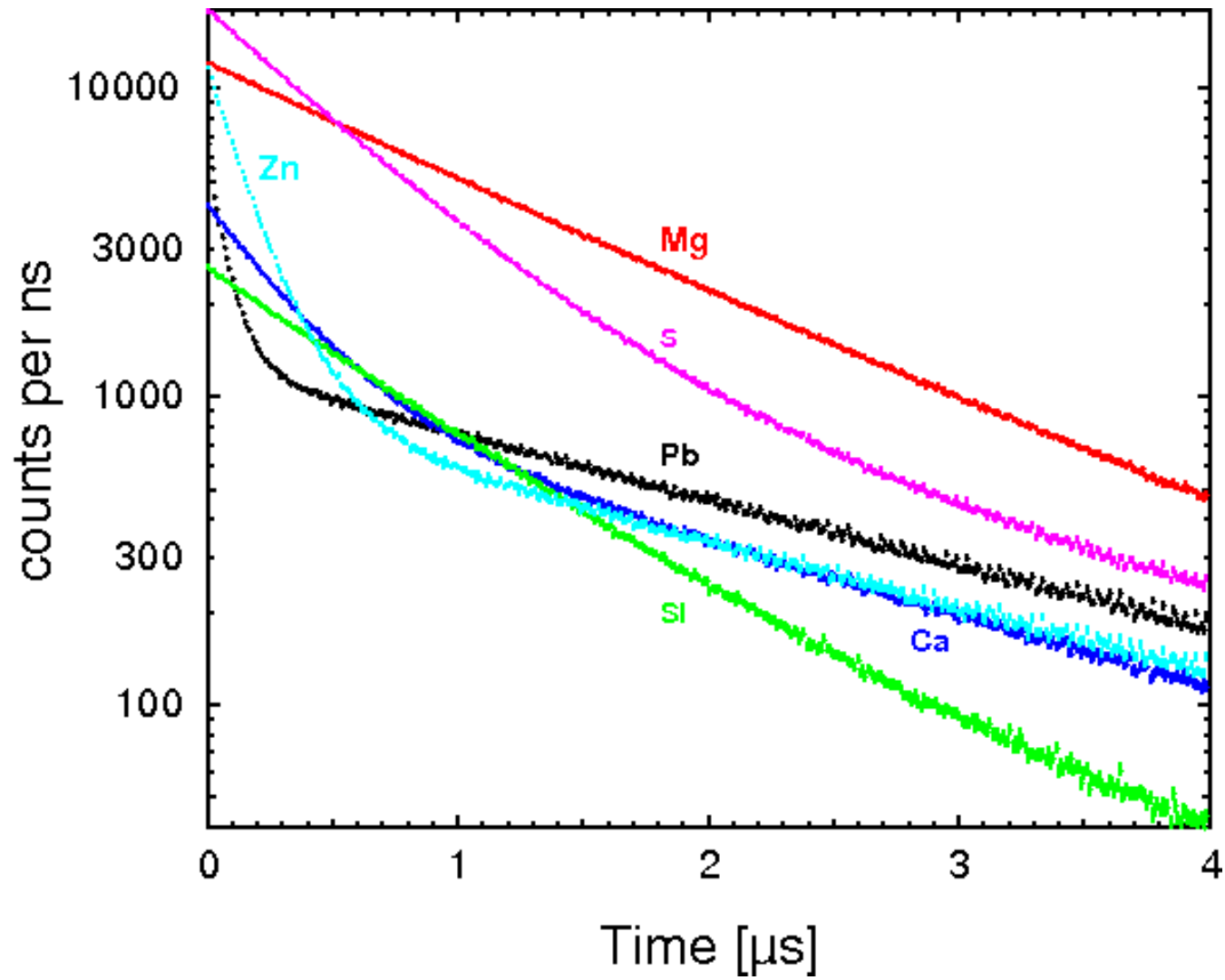
Nuclear μ^- Capture

PROBLEM

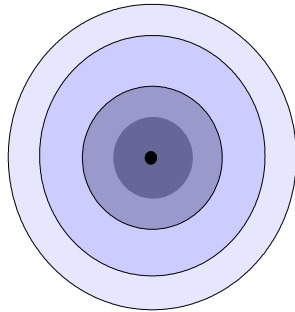
$\mu^- p \rightarrow n \nu_\mu$ in a nucleus:
 rate comparable to that of
 $\mu^- \rightarrow e^- \nu_\mu \bar{\nu}_e$ for $Z \geq 10$.

Possible Help: Many times a **fast neutron** is emitted from nuclear muon capture. Very few measurements have been made of the *correlation* of that neutron with the *muon's spin direction*. If cases are found where this **neutron asymmetry** is sizeable, we *may* be able to do **neutron-triggered μ^-SR** , for which the *event rate* can be **higher** than in **μ^+SR** .

μ^- SR Histograms from E932 (M9B, 2004)



Atomic Capture & $L \cdot S$ Depolarization of μ^-



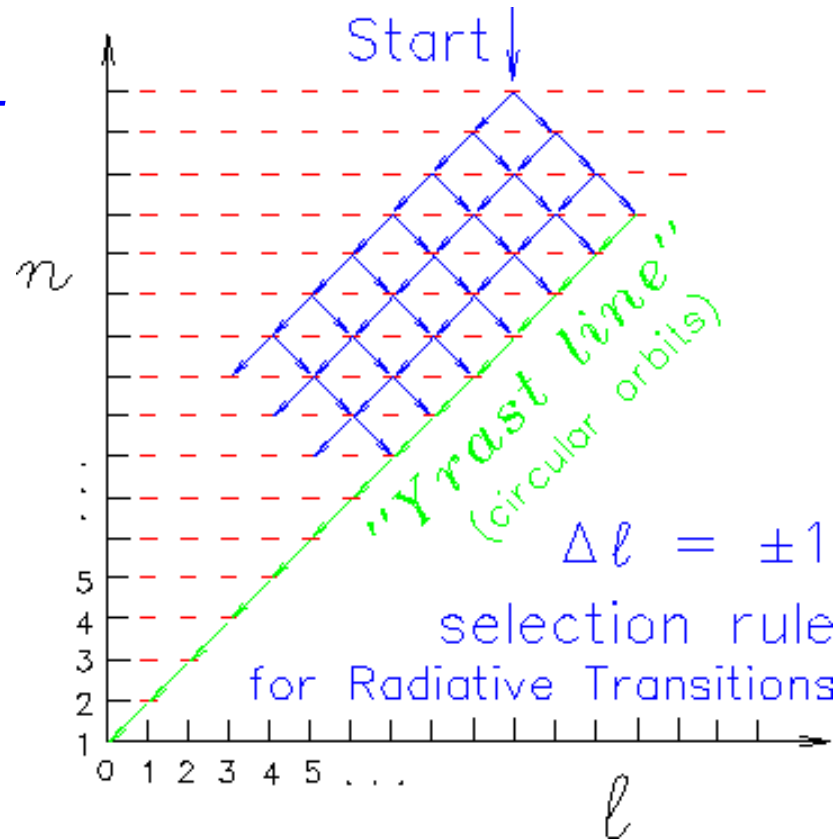
Large impact parameters are more probable
 \Rightarrow initial orbits tend to be fairly **circular**.

View along μ^- momentum

Primitive Atomic Physics:

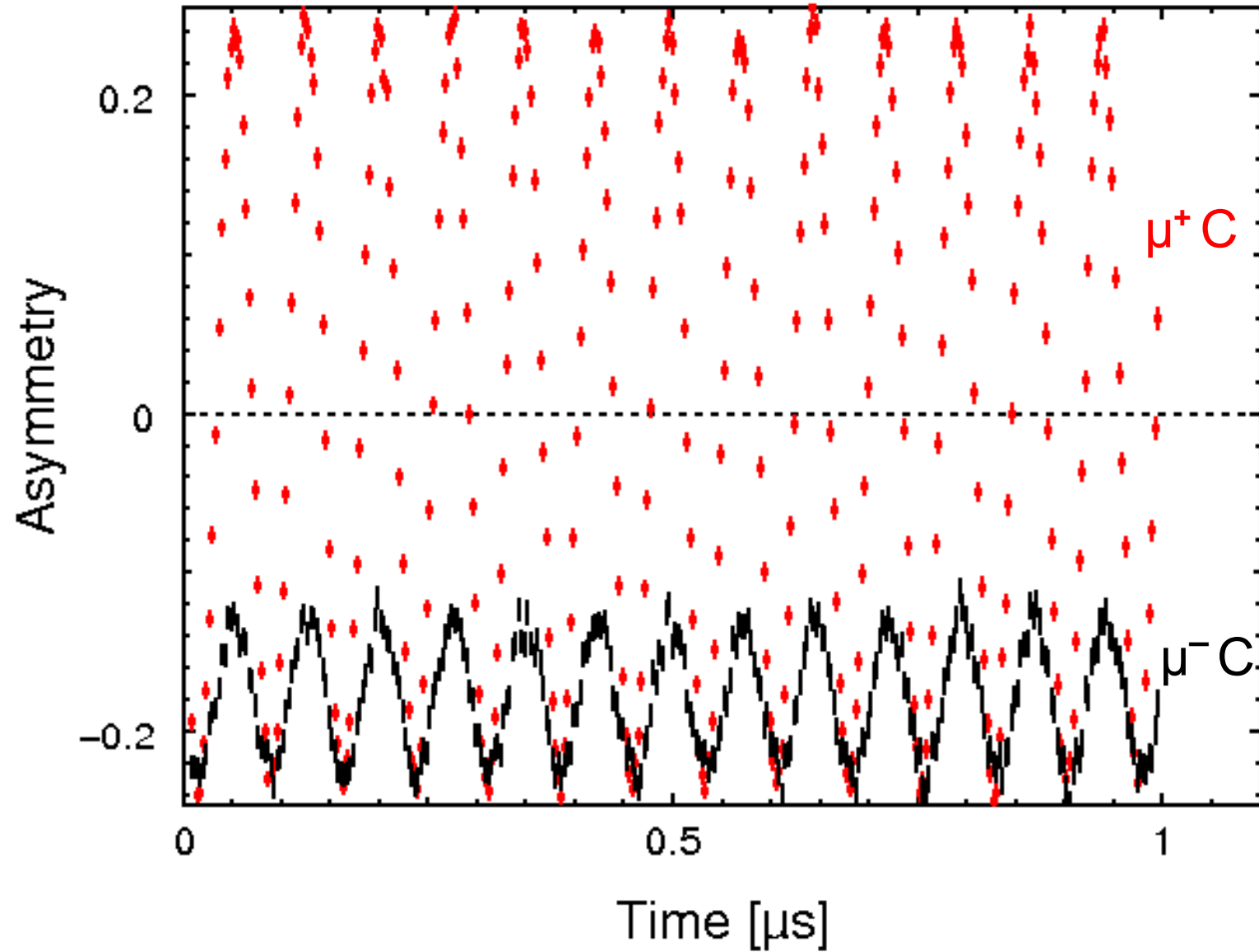
$$r_n = \frac{a_0}{Z} \left(\frac{m_e}{m} \right) n^2$$

$$E_n = - \frac{13.6 \text{ eV}}{n^2} Z^2 \left(\frac{m}{m_e} \right)$$



$L \cdot S$ couplings Depolarize μ^- Spin unless **fast Auger**!

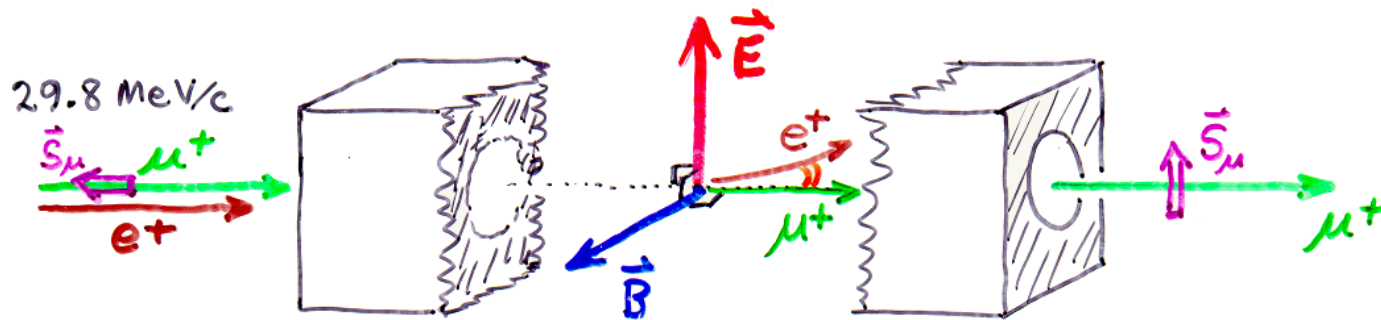
Initial Depolarization of Negative Muons in 1 kG TF



For positive "surface muons"

$E \times B$ velocity selector

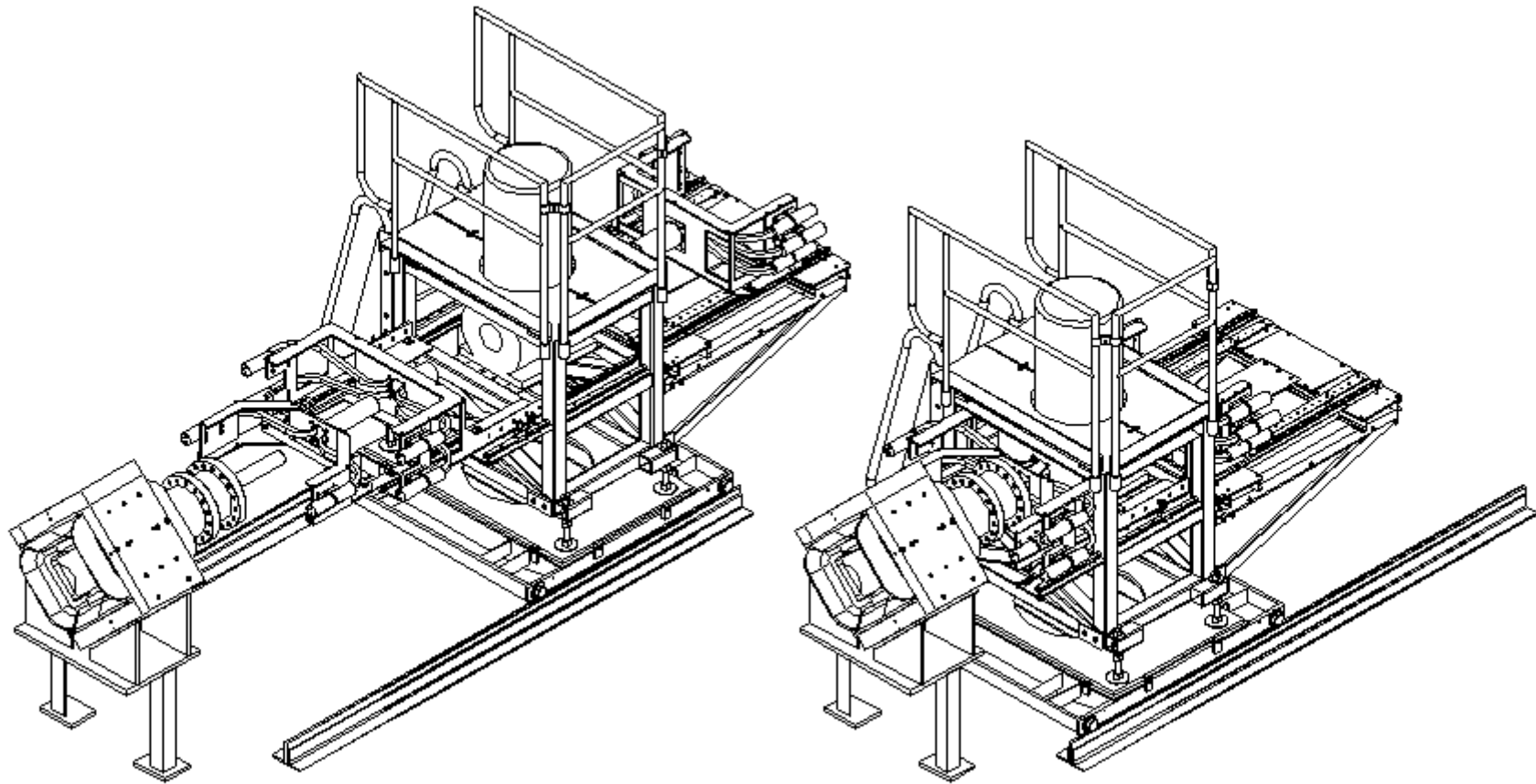
("DC Separator" or Wien filter)



- Removes beam positrons
- Allows TF- μ^+ SR in high field
(otherwise B_{\perp} deflects beam)

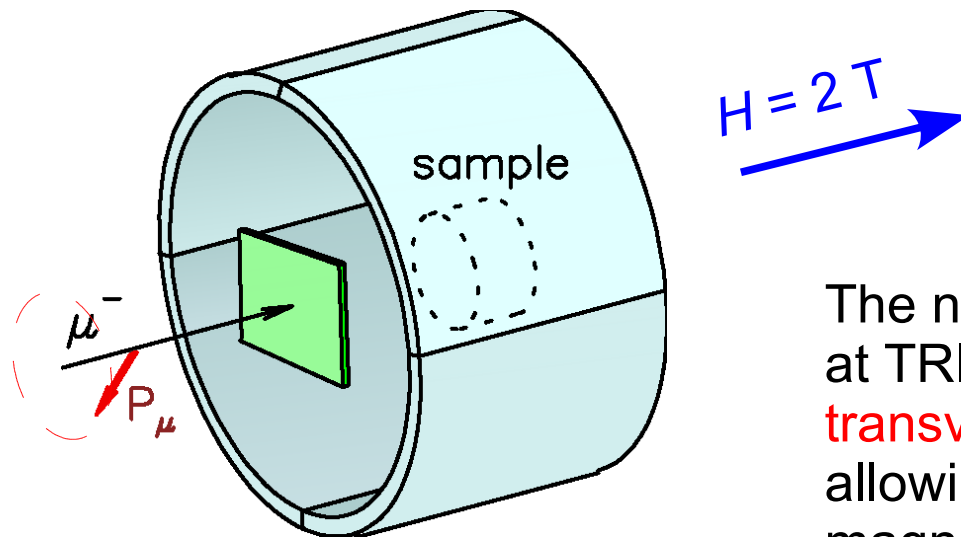
But not available for **negative** muons!

Helios



The *Helios* μ SR spectrometer of the TRIUMF CMMS facility.

Helios on M9B

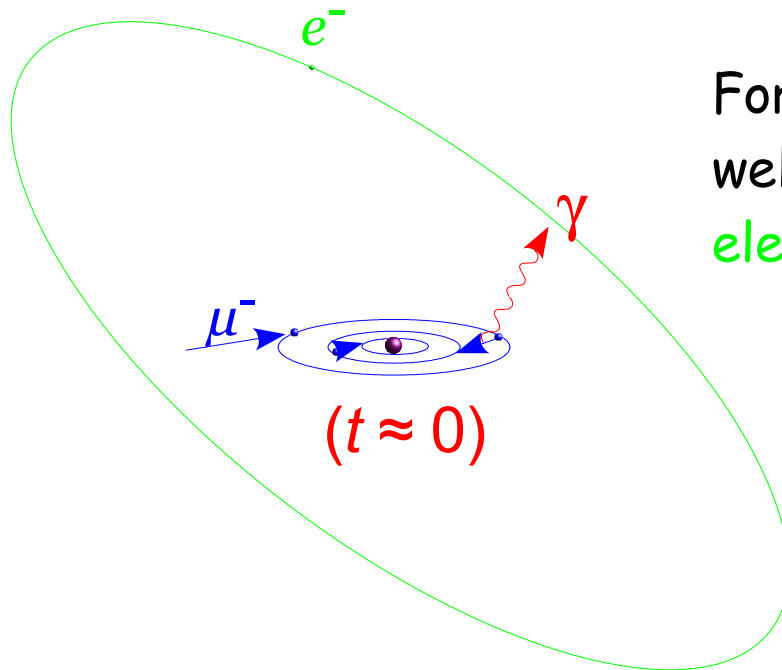


The *Helios* μ SR spectrometer of the TRIUMF CMMS facility enables TF- μ SR at fields up to 2 T, using 4 *e* detectors in a cylindrical array around the target sample.

The negative muon beam of M9B at TRIUMF has nearly 50% **transverse spin polarization**, allowing injection into a strong magnetic field parallel to the beam momentum but (partially) transverse to the spins. Strong TF allows high precision measurements of the muon Larmor frequency and thus of g_μ .

Deeply Bound Hydrogenic States

Muonic orbitals are 207 times smaller than **electronic**.



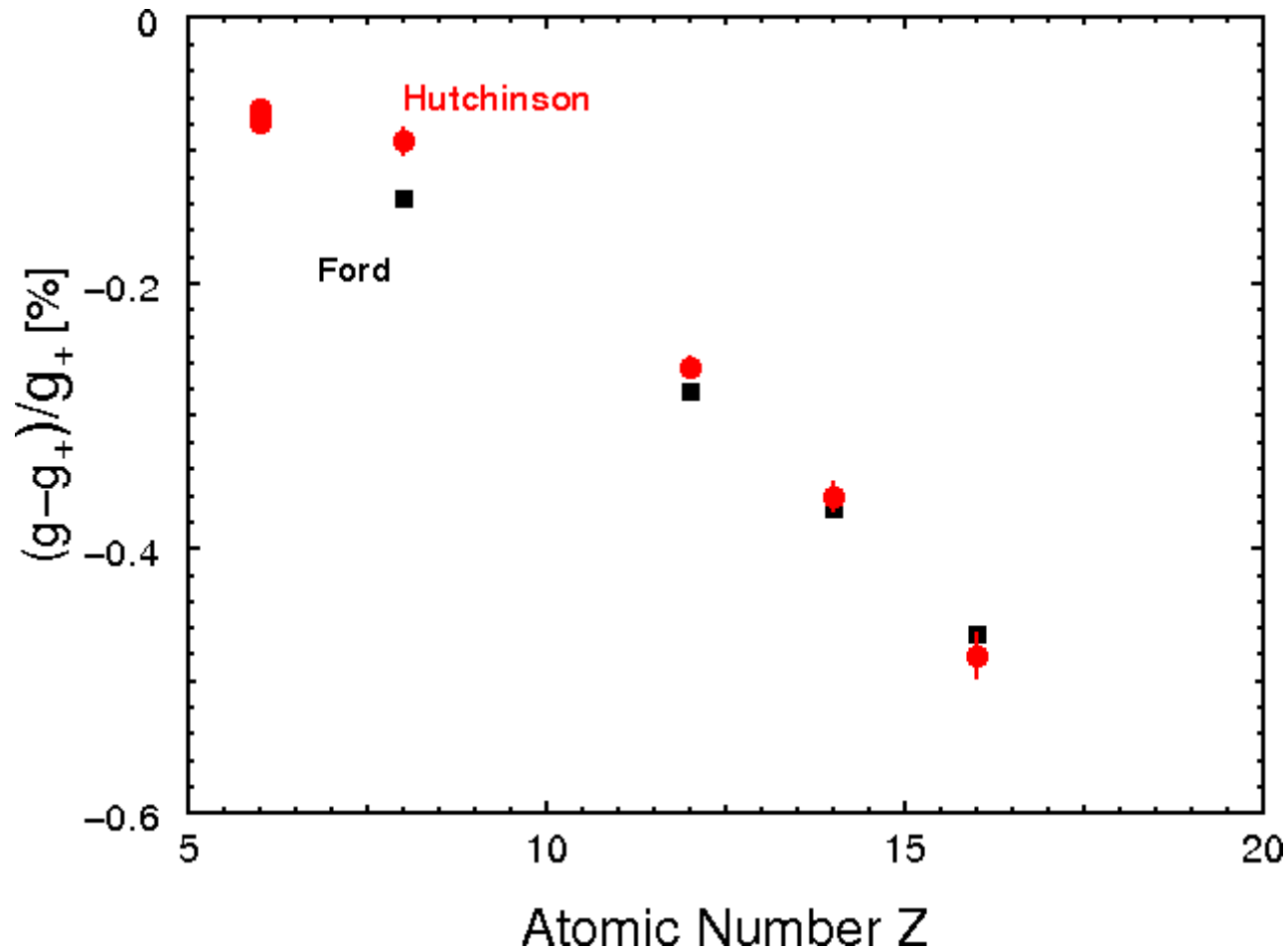
For high Z , the **muon** is well **inside** even core **electronic** orbitals.

Relativistic shift of lepton's magnetogyric ratio (Breit, 1924):
for pointlike nuclei,

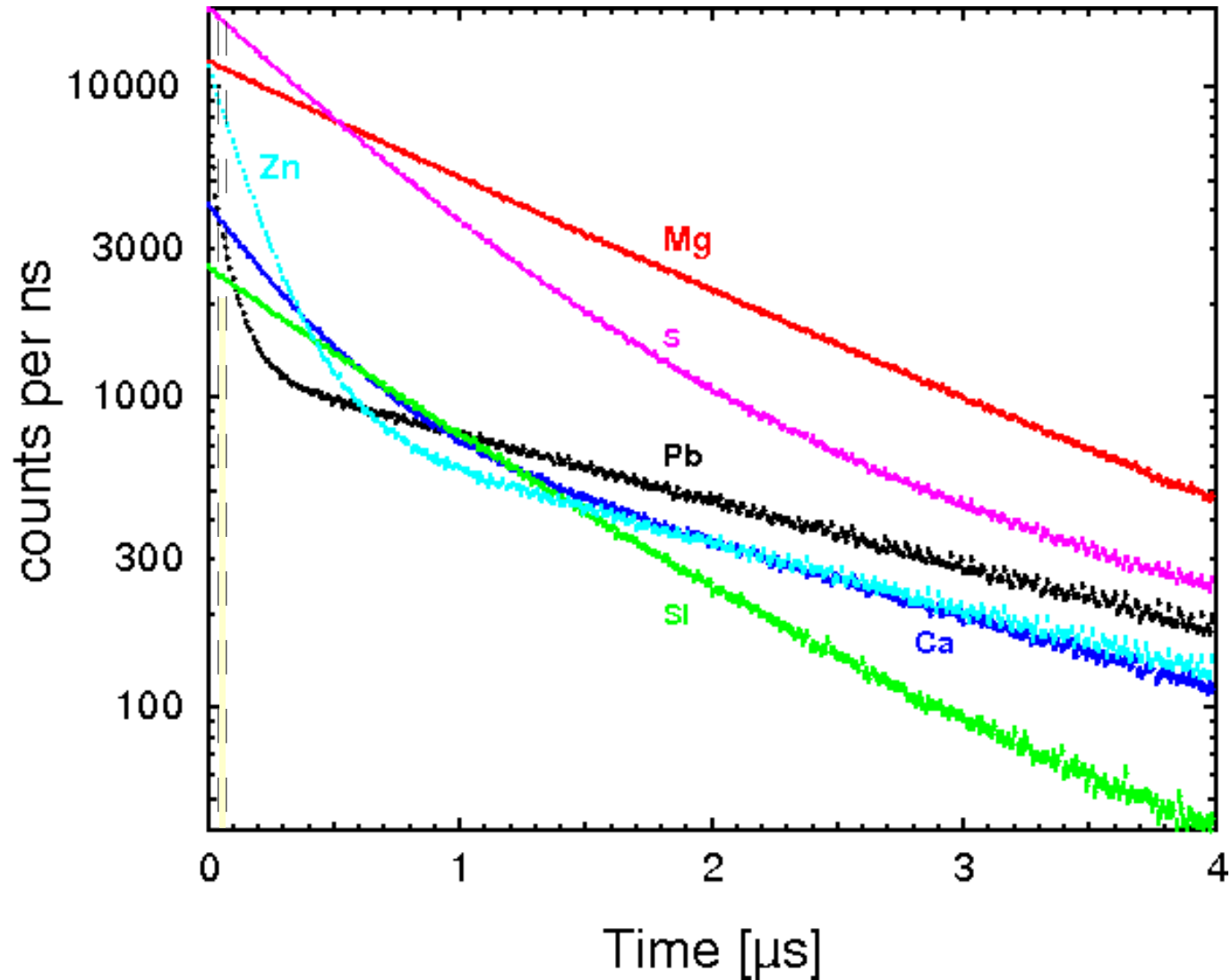
$$\frac{g_{\text{free}} - g}{g_{\text{free}}} = \frac{2}{3} \left(1 - \sqrt{1 - \alpha^2 Z^2} \right) \approx \frac{1}{3} \left(\frac{\bar{v}}{c} \right)^2$$

Theory (Ford et al.) vs. Expt. (Hutchinson et al.) in 1961

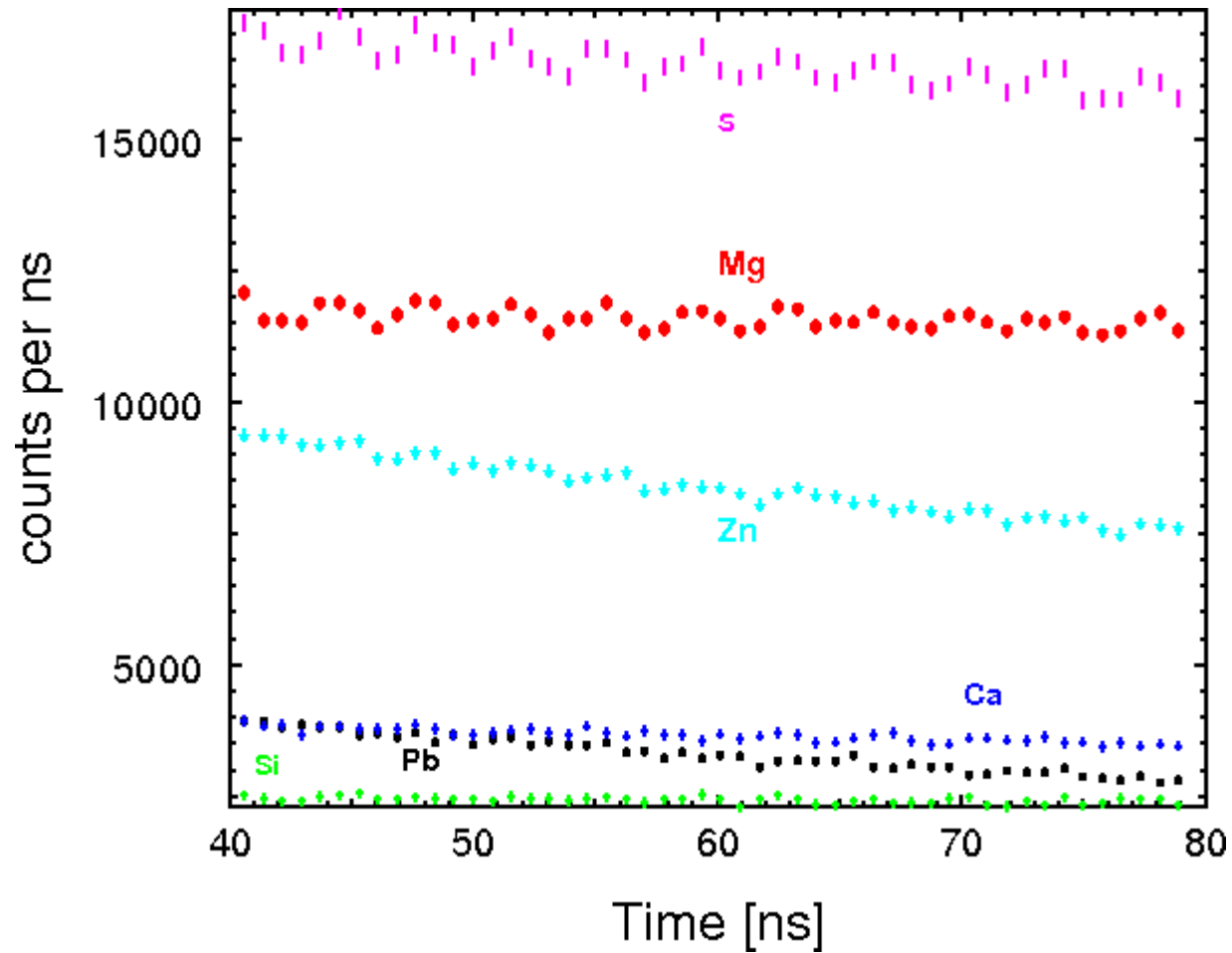
(4 years after the discovery of P violation in π - μ - e decay)



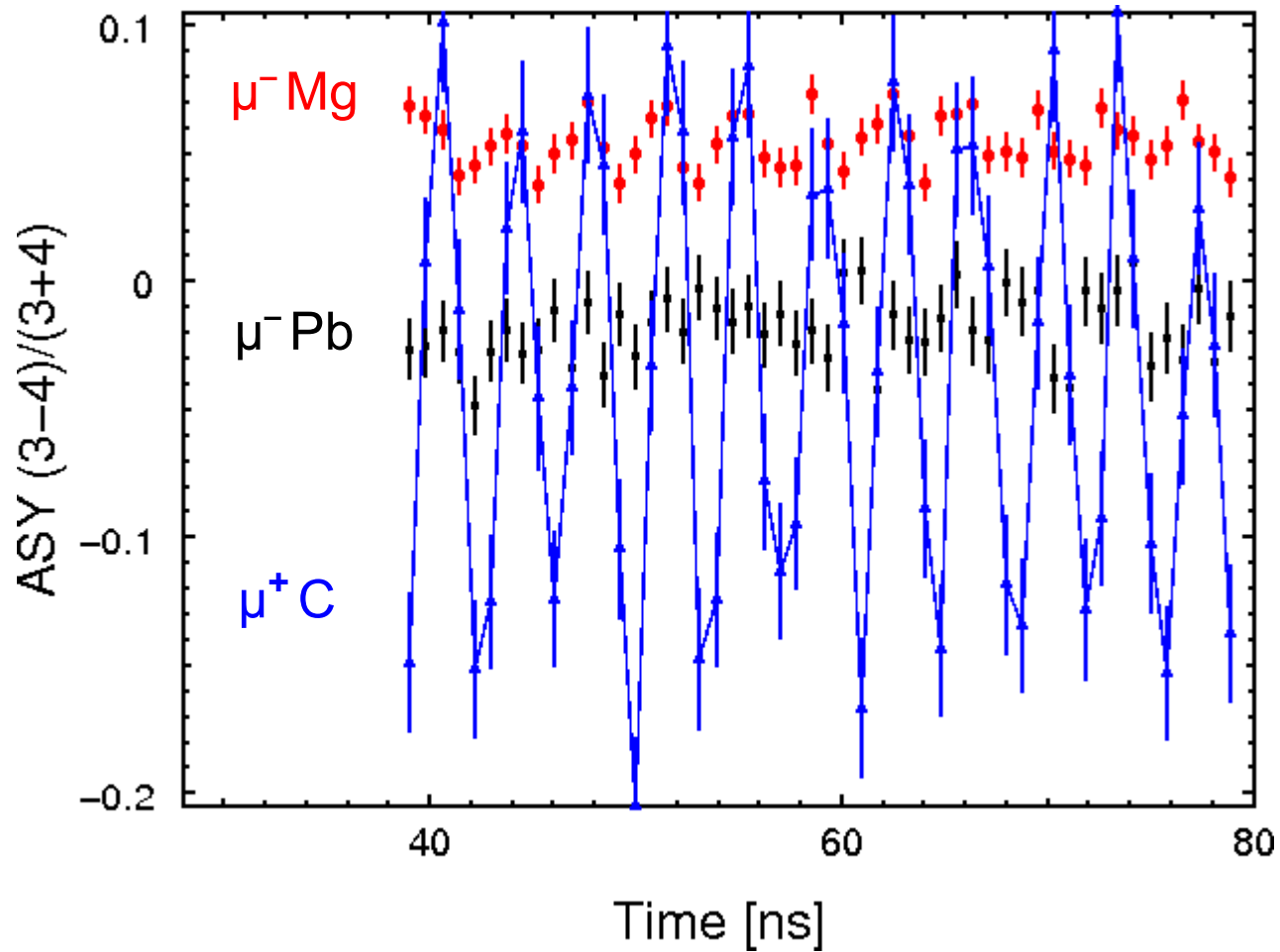
μ^- SR Histograms from E932 (M9B, 2004)



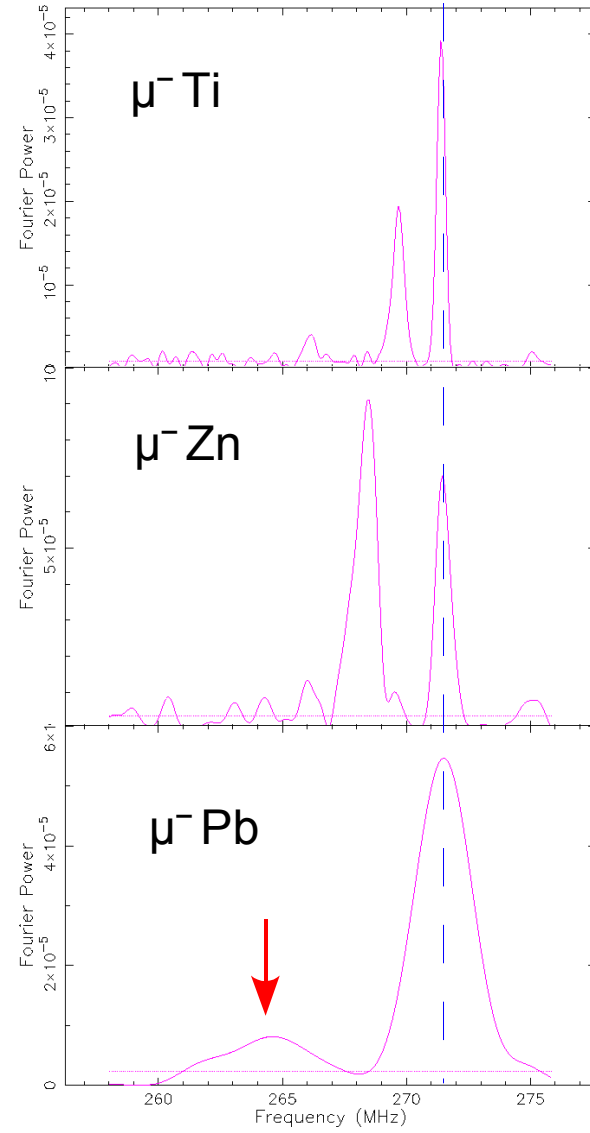
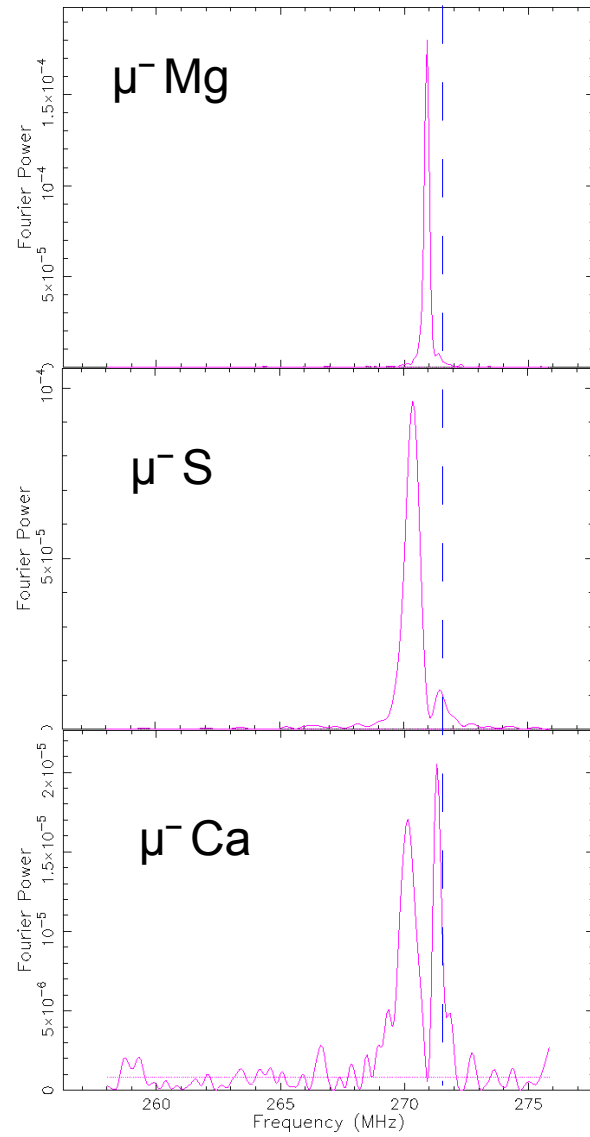
μ^- precession signals from E932 (M9B, 2004)



μ^- precession signals from E932 (M9B, 2004)



μ^- Z frequency spectra from E932 (M9B, 2004)



Raw Data

Sample	Frequency [MHz]
μ^+ in graphite	271.69888 ± 0.00072
μ^+ in Al metal	271.58520 ± 0.00038
μ^- on ^{12}C (graphite)	271.3684 ± 0.0016
μ^- on ^{16}O (H_2O)	271.258 ± 0.010
μ^- on ^{24}Mg (metal)	270.9259 ± 0.0027
μ^- on ^{28}Si	270.6502 ± 0.0069
μ^- on ^{32}S (powder)	270.406 ± 0.008
μ^- on ^{40}Ca (metal)	270.164 ± 0.069
μ^- on Ti (metal)	269.719 ± 0.066
μ^- on Zn (metal)	268.440 ± 0.072
μ^- on Cd (metal)	$265.73^{+0.46}_{-0.57}$
μ^- on Pb (metal)	$264.50^{+0.59}_{-0.62}$

Only *statistical* uncertainties are shown, to emphasize the potential accuracy of such measurements.

In this experiment, *systematic* uncertainties were dominant for the *light* elements.

Results

Sample	g_μ Shift [%]
μ^+ in graphite	0.0499 ± 0.0023
μ^+ in Al metal	0.0080 ± 0.0004
μ^- on ^{12}C (graphite)	-0.0718 ± 0.0023
μ^- on ^{16}O (H_2O)	-0.1124 ± 0.0042
μ^- on ^{24}Mg (metal)	-0.2348 ± 0.0025
μ^- on ^{28}Si	-0.3363 ± 0.0034
μ^- on ^{32}S (powder)	-0.4262 ± 0.0036
μ^- on ^{40}Ca (metal)	-0.5155 ± 0.025
μ^- on Ti (metal)	-0.679 ± 0.024
μ^- on Zn (metal)	-1.150 ± 0.026
μ^- on Cd (metal)	$-2.15^{+0.17}_{-0.21}$
μ^- on Pb (metal)	$-2.60^{+0.22}_{-0.23}$

Fractional shifts (in %) of the negative muon's g factor due to ***relativistic*** effects in the deeply bound ground state of the muonic atom.

(In Pb, most of the muon's orbital lies *inside* the nucleus!)

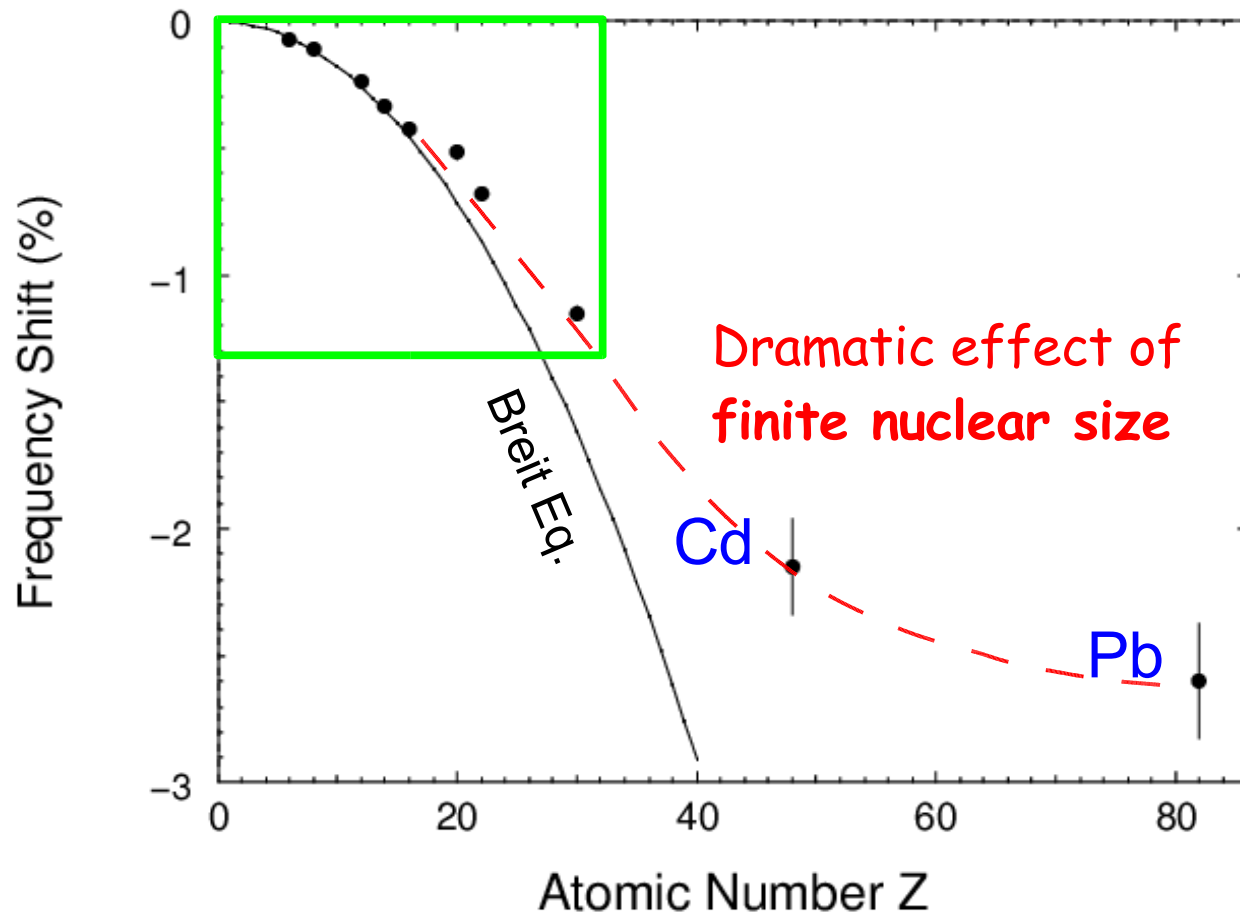
So what does it mean?

For pointlike nuclei (Breit, 1928):

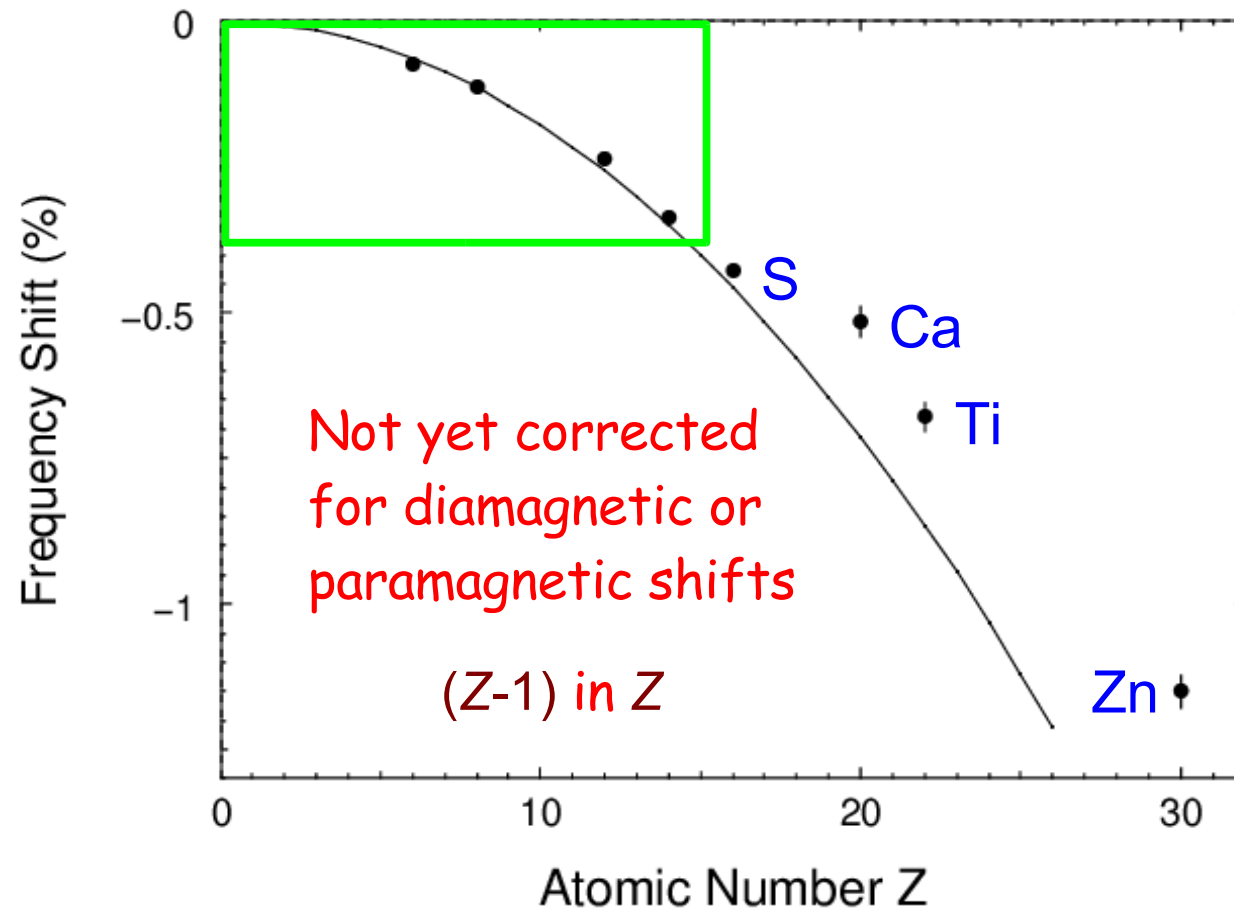
$$\frac{g_{\text{free}} - g}{g_{\text{free}}} = \frac{2}{3} \left(1 - \sqrt{1 - \alpha^2 Z^2} \right) \approx \frac{1}{3} \left(\frac{\bar{v}}{c} \right)^2$$

Improved by Margeneau (1940) and later by Ford *et al.* (1962) in response to first μ^- SR measurements by Hutchinson *et al.* (1961) in light elements. First high-Z measurements by Yamazaki *et al.* (1974) challenged by Mamedov *et al.* (2003). Meanwhile electronic spectroscopy of high Z hydrogenlike ions has become possible [e.g. Häffner *et al.* (2000)].

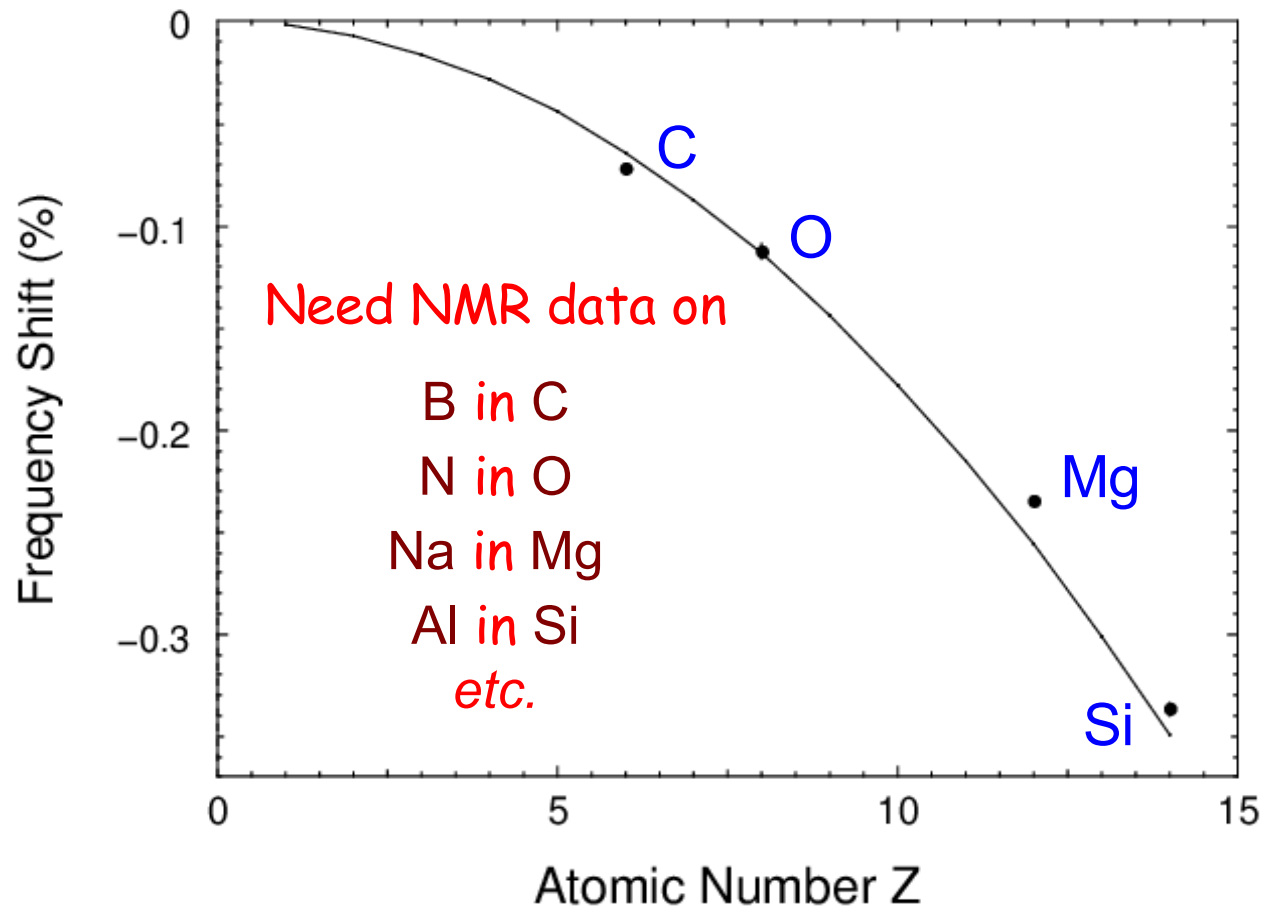
Relativistic Shift of μ^- Frequency



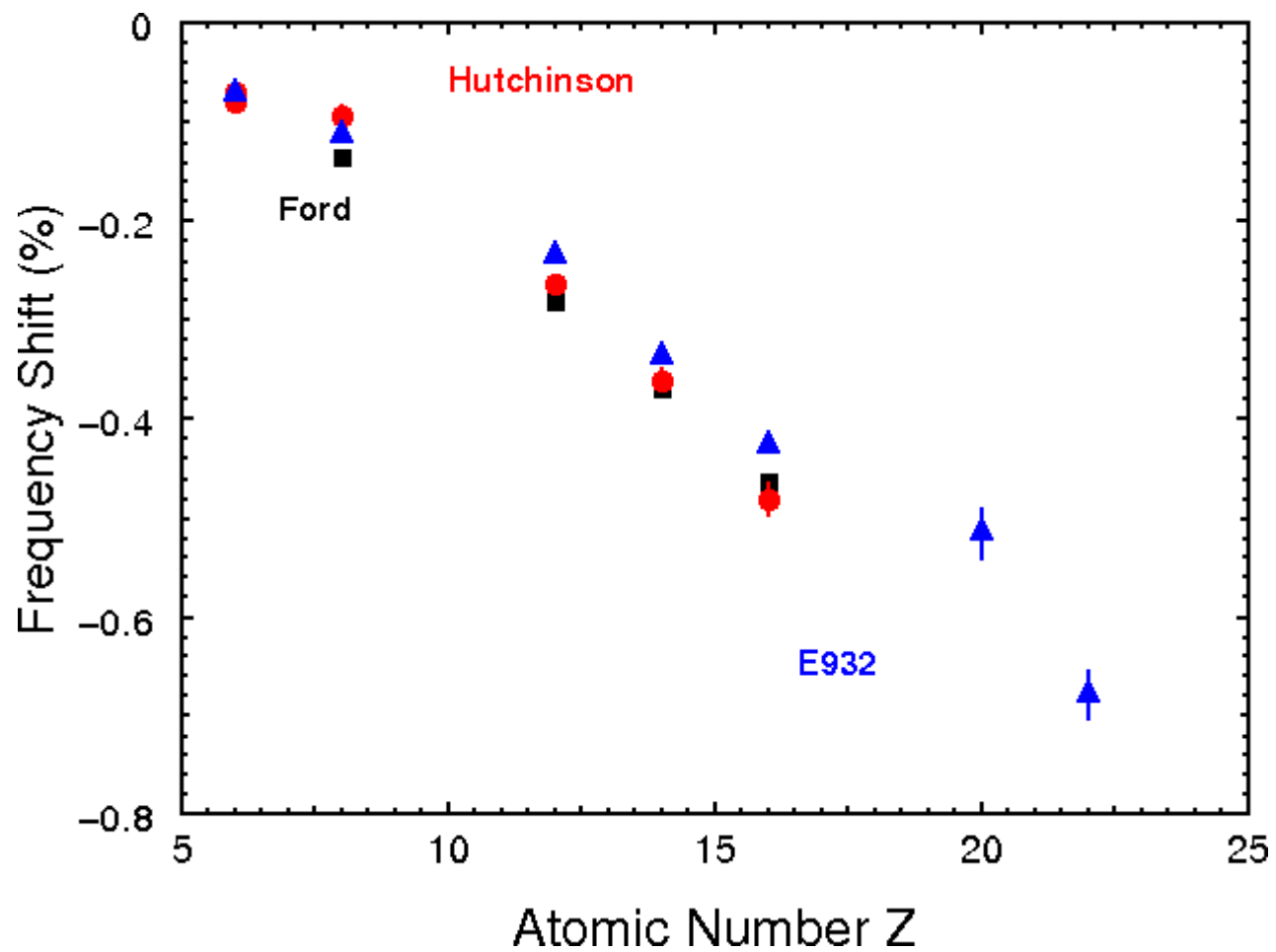
Relativistic Shift of μ^- Frequency



Relativistic Shift of μ^- Frequency



E932 vs. Ford & Hutchinson



Phil Anderson:

(at a High T_c Superconductivity conference)

“Experimentalists should not attempt to interpret their own data.”

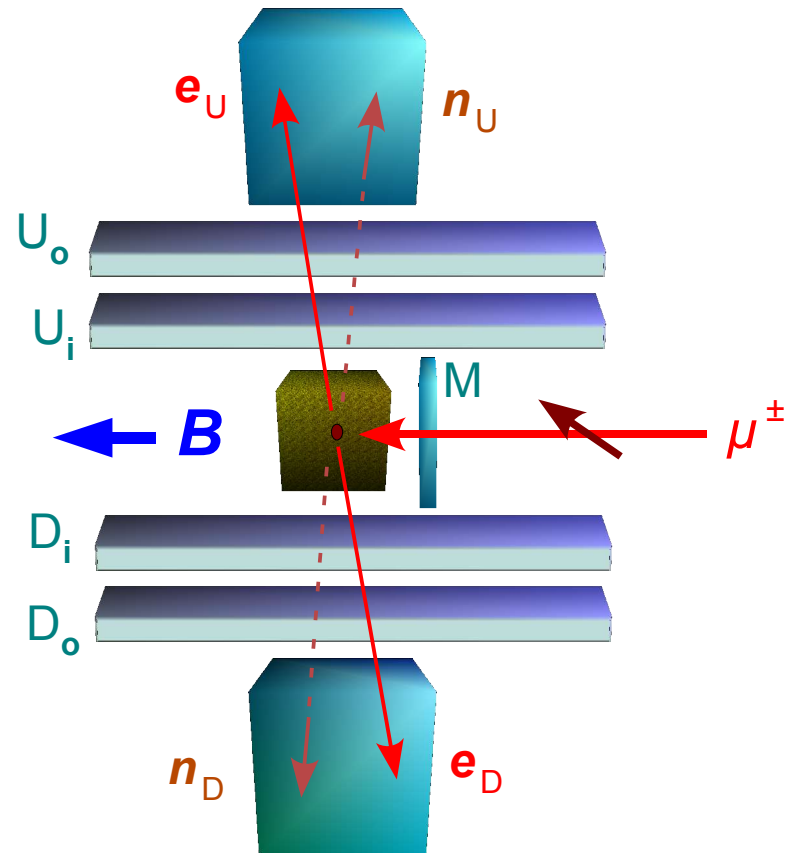
[paraphrased]

Darth Vader:

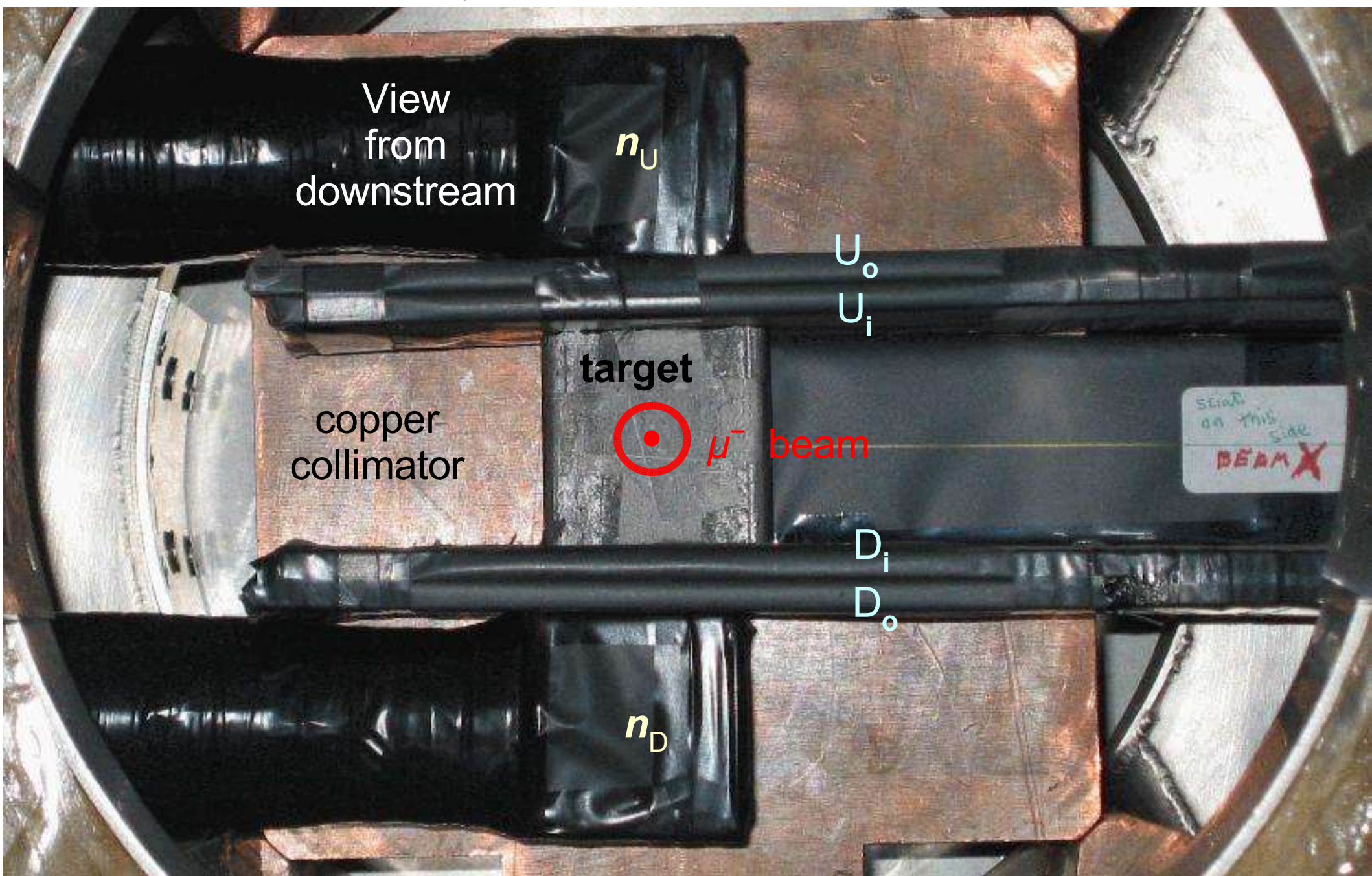
“Leave that to me.”

Finis
(for now)

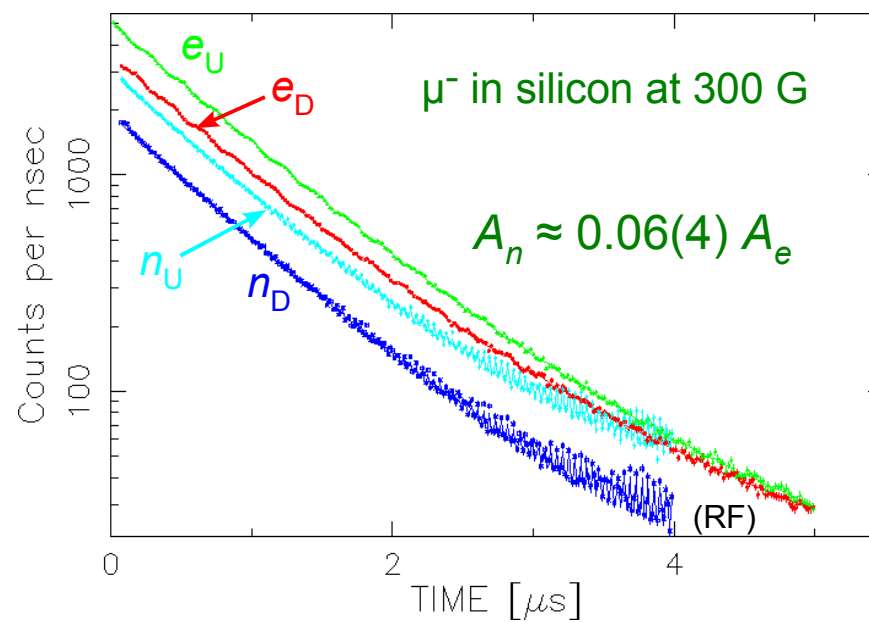
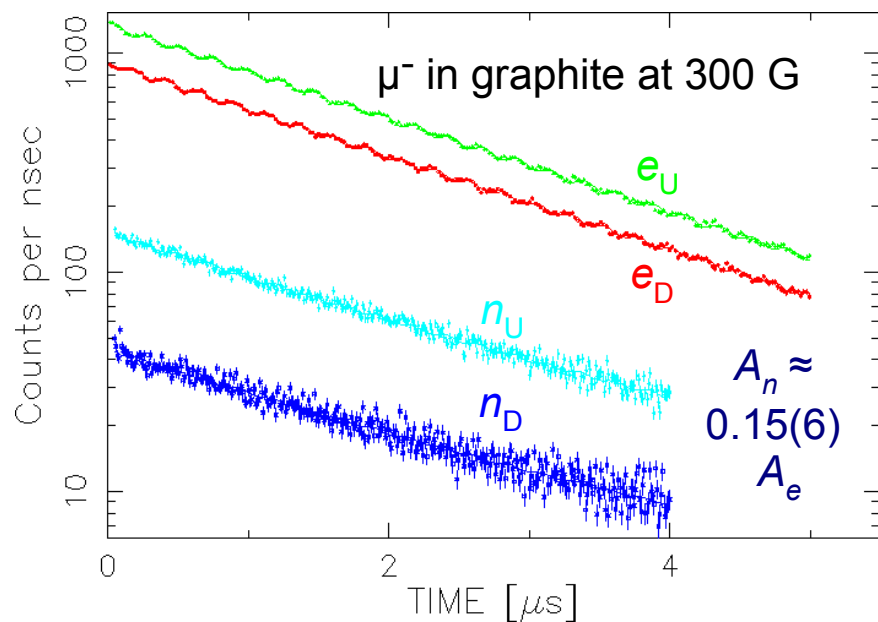
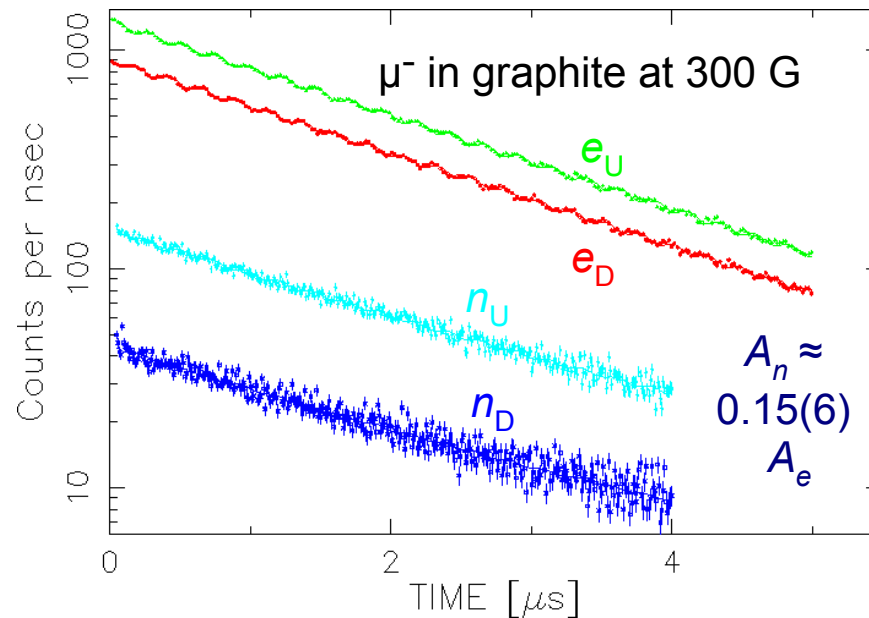
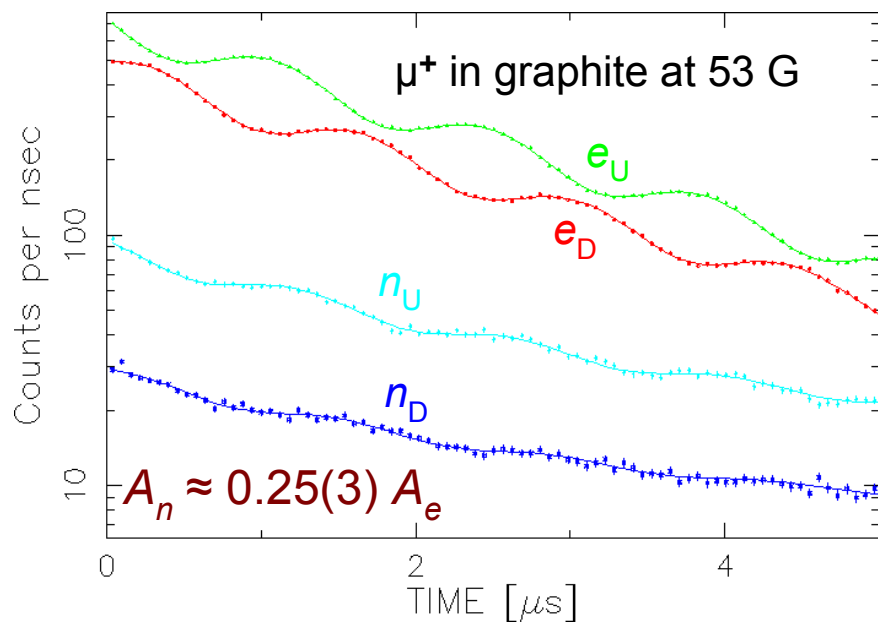
E932 on M9B, July 2005: 1st look at Neutrons



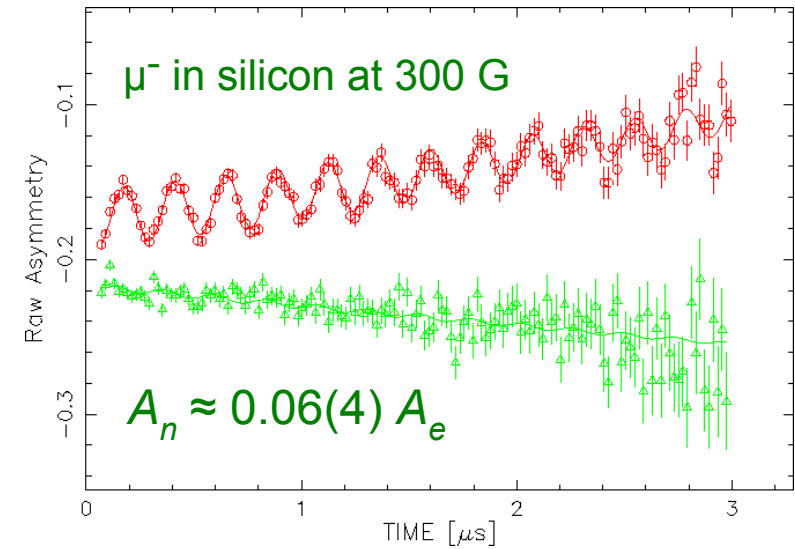
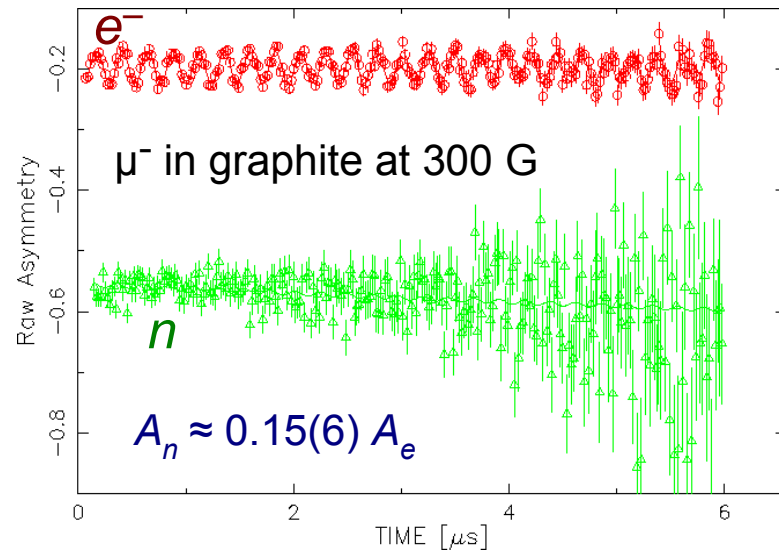
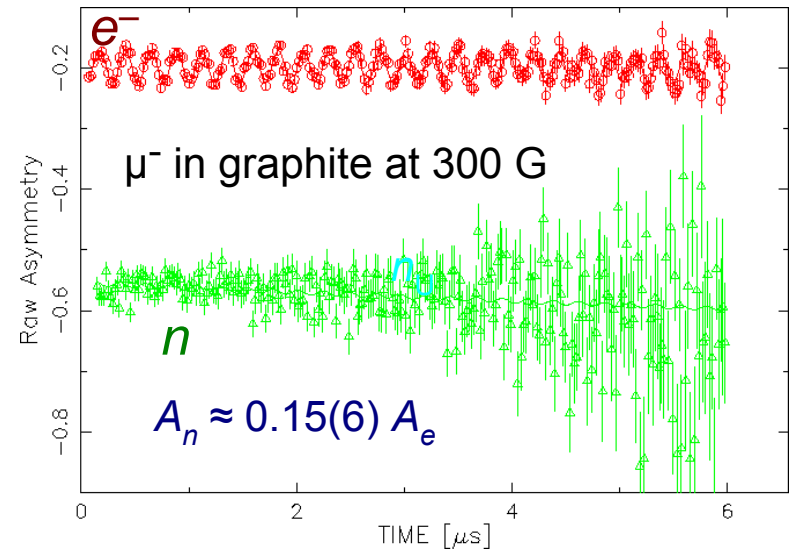
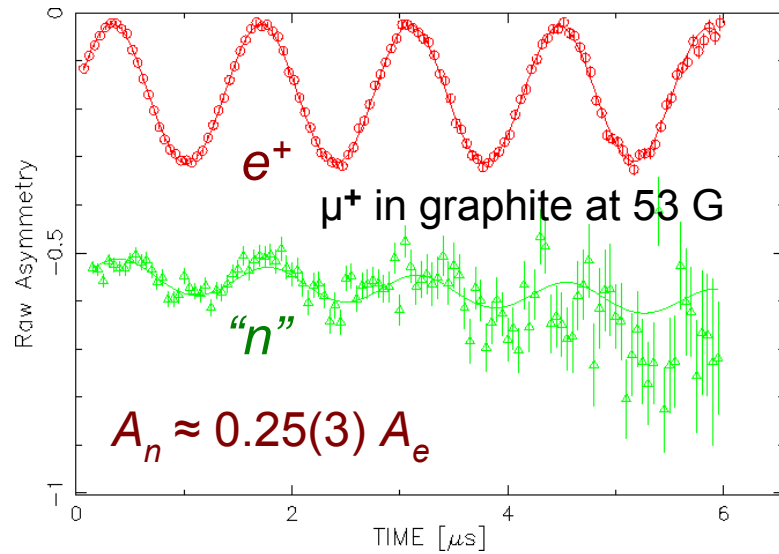
E932 on M9B, July 2005: **Neutron counters** in Omni'



M9B data 2-4 July 2005: Raw Histograms



M9B data 2-4 July 2005: (U-D)/(U+D) Asymmetries



Negative muon spin precession measurement of the hyperfine states of muonic sodium

Jess H. Brewer, Khashayar Ghandi,
Aaron M. Froese & Bruce A. Fryer

CIAR, UBC, TRIUMF & Chatelech Secondary School

Both hyperfine states of muonic ^{23}Na and the rate R_{HF} of conversion between them have been observed directly in a high field negative muon spin precession experiment using a backward muon beam with transverse spin polarization. The result in metallic sodium, $R = 13.7 \pm 2.2 \mu\text{s}^{-1}$, is consistent with Winston's prediction in 1963 based on Auger emission of core electrons, and with the measurements of Goringe *et al.* in Na metal, but not with their smaller result in NaF. In NaOH we find $R = 23.5 \pm 8 \mu\text{s}^{-1}$, leaving medium-dependent effects ambiguous.

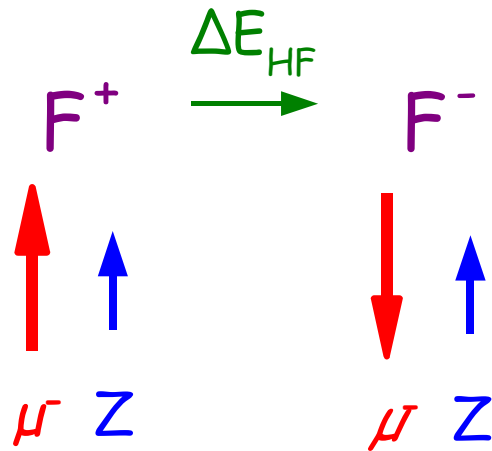
Phys. Rev. C Brief Reports 71, 058501 (2005).

HFS OF SELECTED "BARE" MUONIC ATOMS:

$\mu^-(Z, A)$ ground state $\left\{ \begin{array}{l} F^+ = I + 1/2 \\ F^- = I - 1/2 \end{array} \right.$
 ΔE_{HF}

ELEMENT $\begin{smallmatrix} A \\ Z \end{smallmatrix}$	ΔE_{HF} (eV)	effective field B_{HF} (G)	Lowest e^- SHELL ejected [B.E. (eV)]	$F^+ \rightarrow F^-$ RATE R (μs^{-1})		REF
				THEOR	EXPT	
${}^1\text{H}_1$	0.1817	3.24×10^9	NONE	1.3×10^4	(ρ/ρ_0)	Ponomarev '78 smilga & Fil'chenko '83
${}^3\text{He}_2$	1.373	2.45×10^{10}	NONE	nil		
${}^{11}\text{B}_5$	18	3.2×10^{11}	L_1 [9.3]	0.25	0.33(5)	Winston '63 Favart et al '70
${}^{13}\text{C}_6$	11	2.0×10^{11}	L_2 [8.3]		0.016(12)	BOAM '83
${}^{14}\text{N}_7$	7.5	1.3×10^{11}	NONE	nil	0.092(33)	"
${}^{19}\text{F}_9$	126	2.2×10^{12}	L_1 [30]	5.8	6.1(7) (neutrons)	Winston '63
${}^{27}\text{Al}_{13}$	263	4.7×10^{12}	L_1 [89]	41	41(9)	" & Brewer '83
${}^{51}\text{V}_{23}$	1220	2.1×10^{13}	L_1 [565]	700		Winston '63
${}^{93}\text{Nb}_{41}$	~ 5000	$\sim 10^{14}$				
${}^{209}\text{Bi}_{83}$	4660	8.3×10^{13}	MEASURED VIA SPLITTING OF X-RAY ENERGY			

Characteristic
precession frequencies
of F^\pm hyperfine states
in selected low- Z
muonic atoms

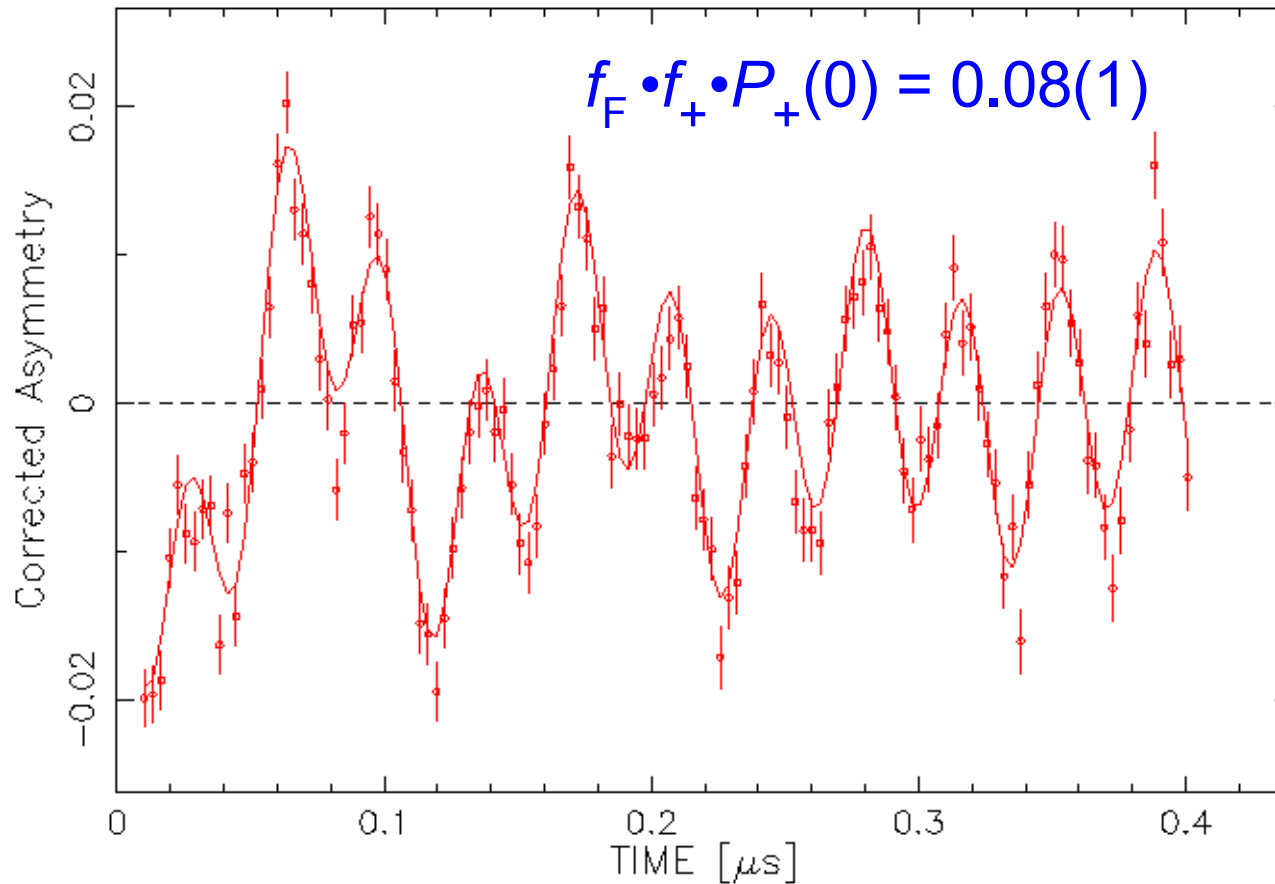


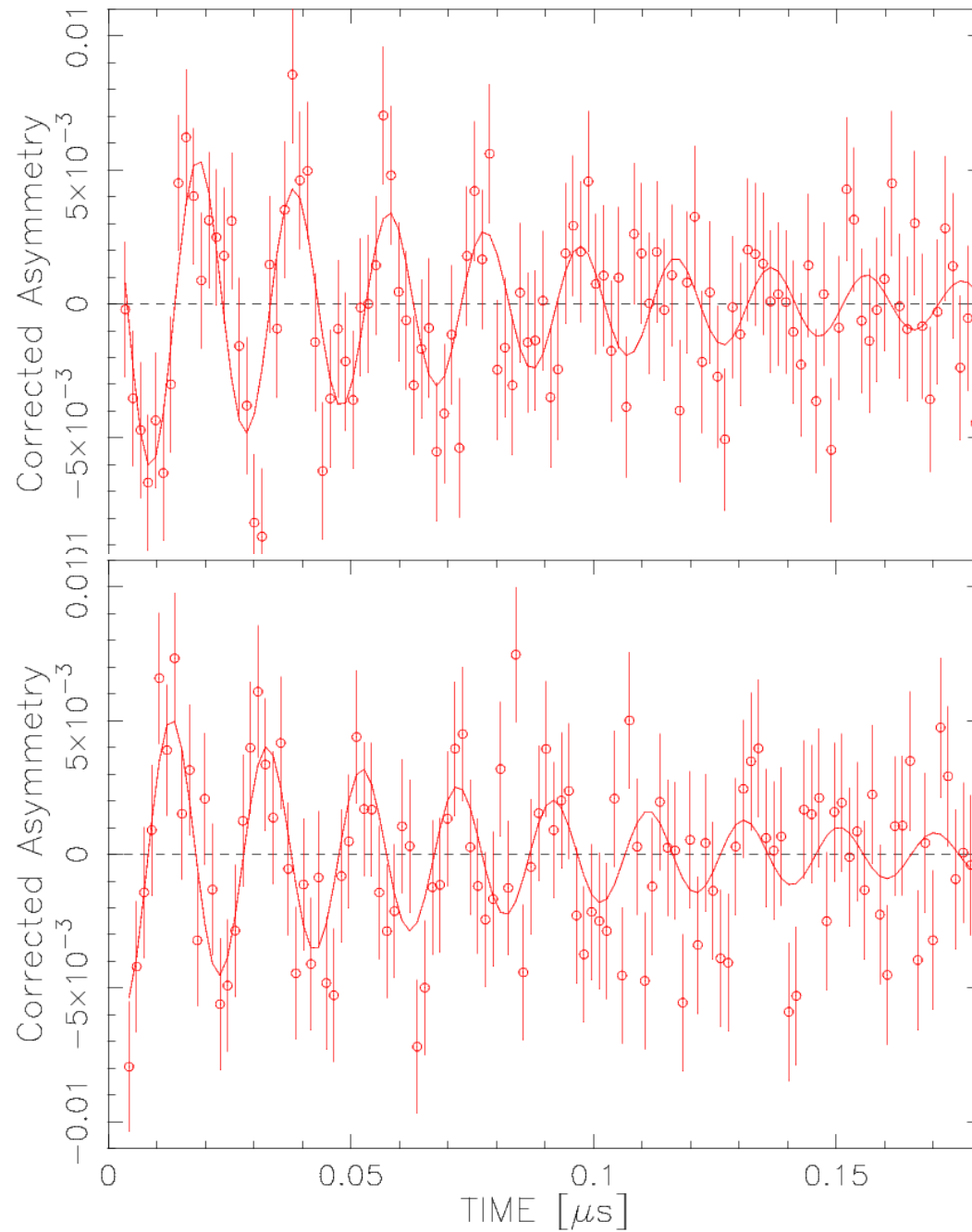
Isotope ${}^A E I_Z$	Nucl. Spin	Natural Abundance	Moment μ_N/μ_μ	Frequency Ratios	
				F^+/μ^-	F^-/μ^-
${}^1\text{H}_1$	1/2	≈ 1	-0.314109	0.342946	0
${}^2\text{H}_1$	1	≈ 0	-0.096436	0.301188	-0.397624
${}^6\text{Li}_3$	1	0.07	-0.092454	0.302515	-0.394969
${}^7\text{Li}_3$	3/2	0.93	-0.366253	0.158437	-0.402606
${}^9\text{Be}_4$	3/2	≈ 1	0.132447	0.283112	-0.194814
${}^{10}\text{B}_5$	3	0.19	-0.202528	0.113925	-0.181434
${}^{11}\text{B}_5$	3/2	0.81	-0.302380	0.174405	-0.375992
${}^{13}\text{C}_6$	1/2	0.01	-0.079000	0.460500	0
${}^{14}\text{N}_7$	1	≈ 1	-0.045394	0.318202	-0.363596
${}^{19}\text{F}_9$	1/2	≈ 1	-0.295666	0.352167	0
${}^{23}\text{Na}_{11}$	3/2	≈ 1	-0.249406	0.187648	-0.353919
${}^{25}\text{Mg}_{12}$	5/2	0.10	0.096197	0.182700	-0.144221
${}^{27}\text{Al}_{13}$	5/2	≈ 1	-0.409555	0.098408	-0.262229

μ^-SR in Teflon $[(CF_2)_n]$: High frequency signal: $^{12}C \mu^-$

Low frequency signal:

F^+ (triplet) state of $^{19}F\mu^-$ $\xrightarrow[\text{HF transition rate}]{R_{HF} = 5.2(5) \mu s^{-1}}$ F^- (singlet) state

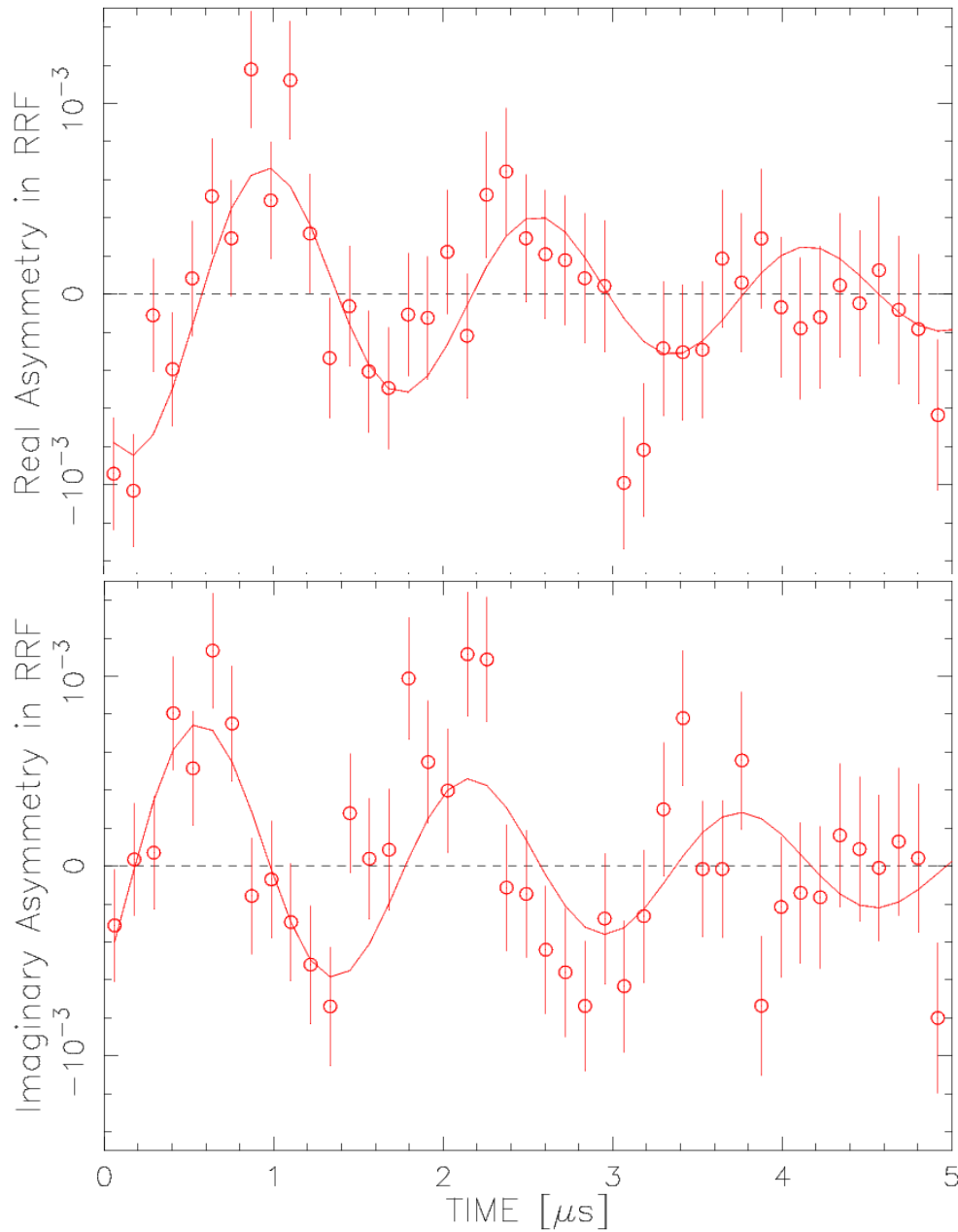




$^{23}\text{Na } \mu^-$
 F^+ state

$$f_+ \cdot P_+(0) = 0.064(8)$$

$$R_{\text{HF}} = 14(3) \mu\text{s}^{-1}$$



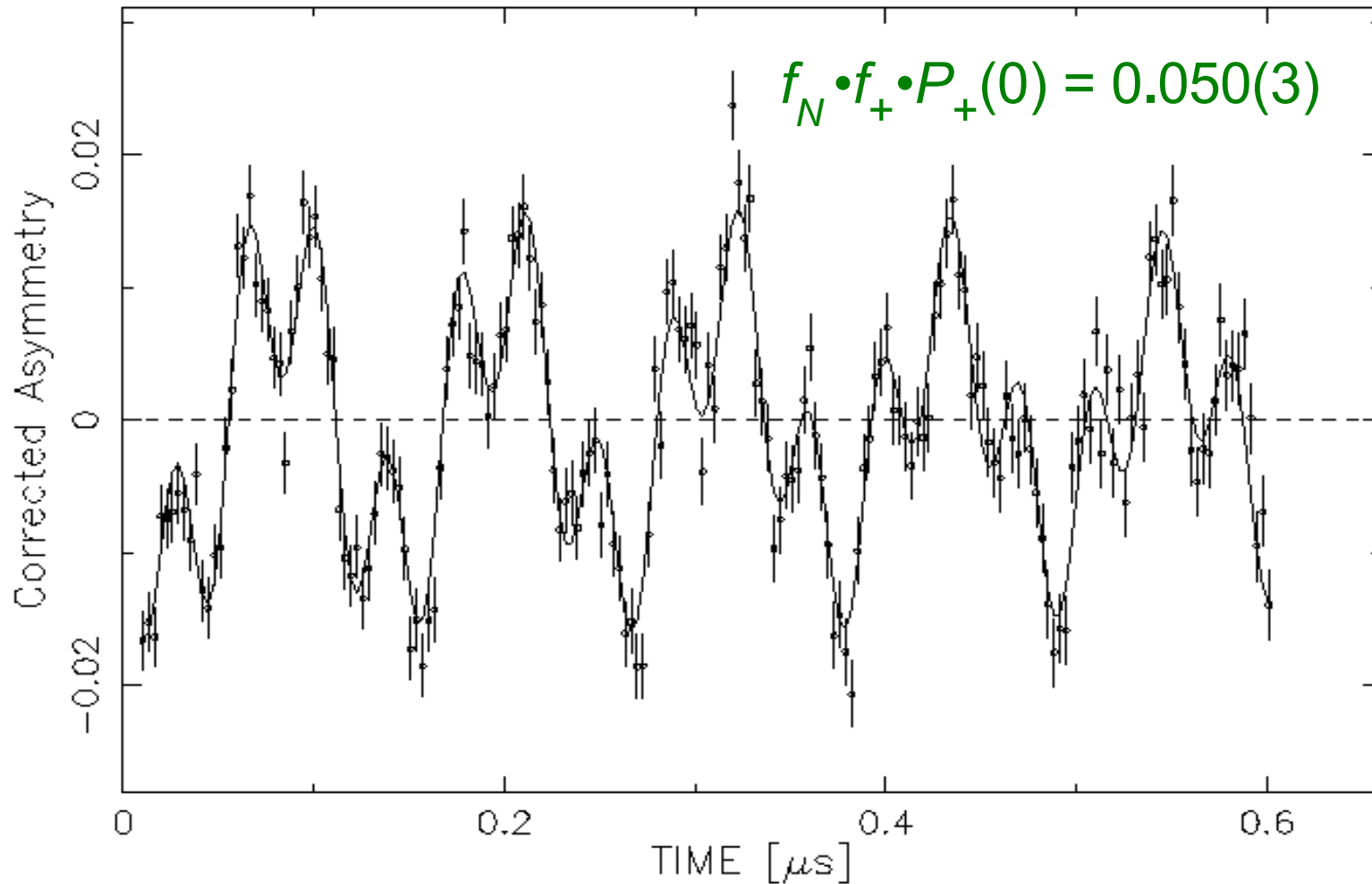
$^{23}\text{Na } \mu^-$
 F^- state

$$f_{-} \cdot P_{-}(0) = 0.0090(15)$$

$$T_2^{-1} = 0.024(100) \mu\text{s}^{-1}$$

μ^-SR in Melamine ($C_3H_6N_6$):

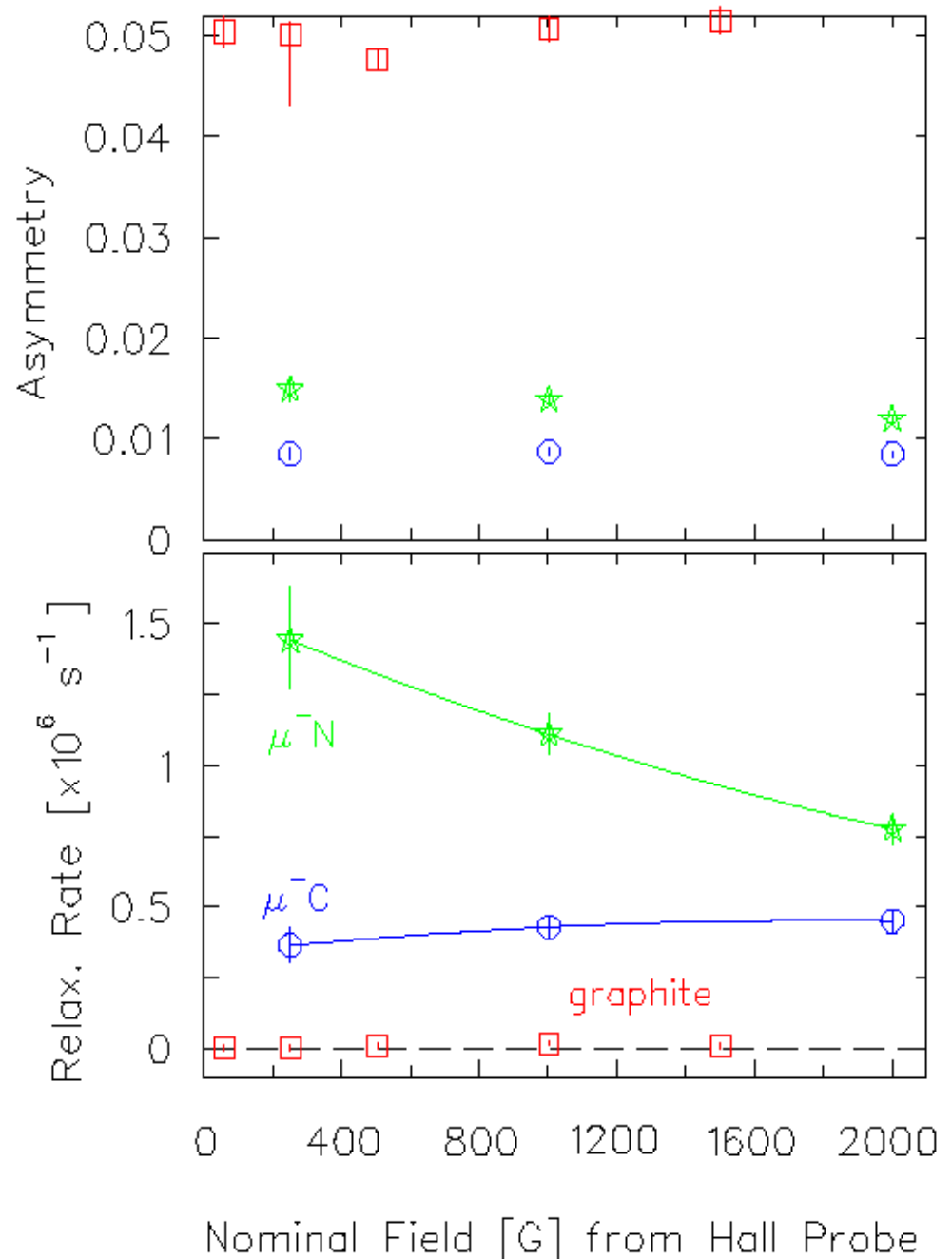
High frequency signal: $^{12}C \mu^-$ Low frequency signal: F^+ state of $^{14}N \mu^-$



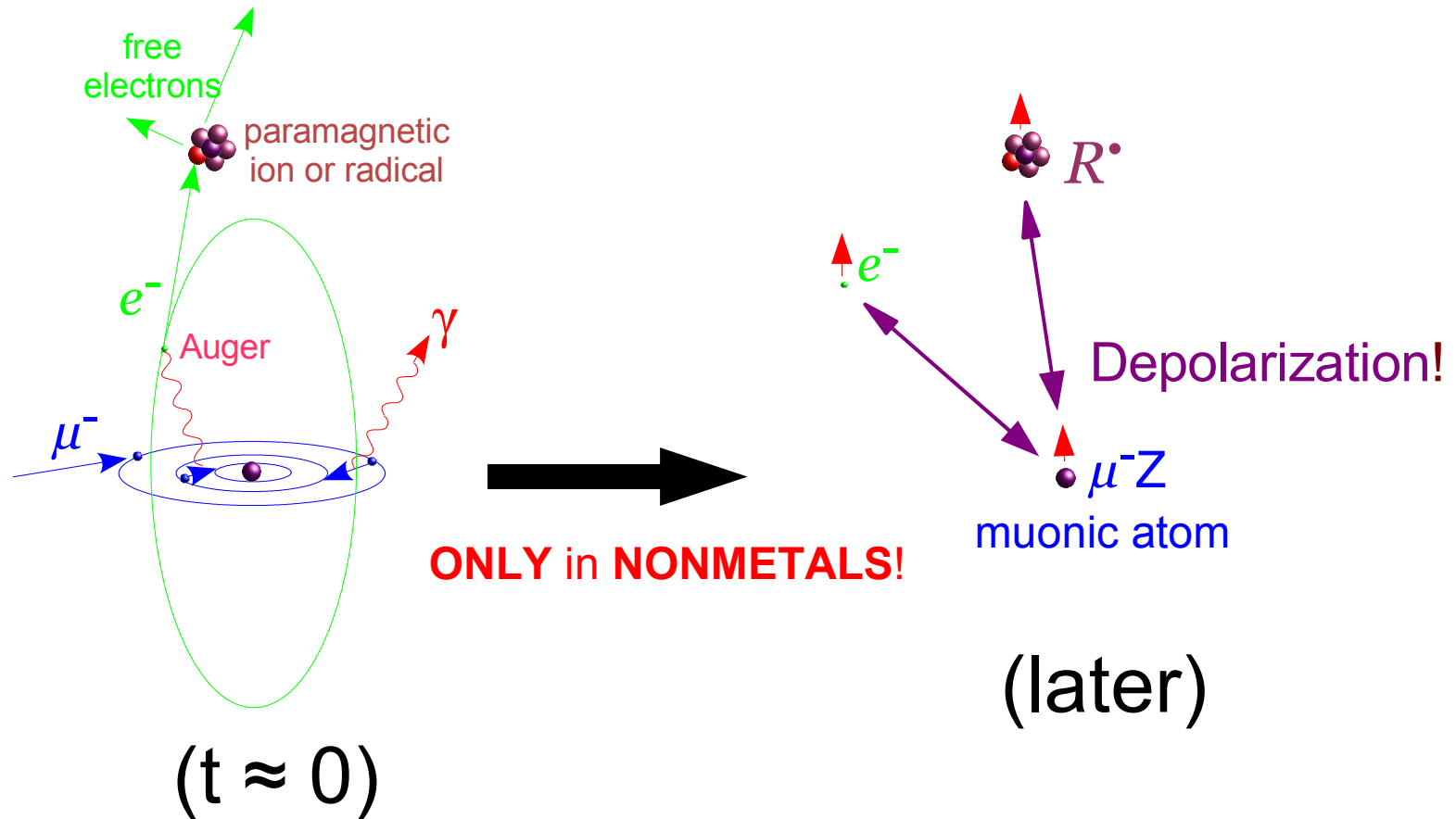
μ^-SR in Melamine ($C_3H_6N_6$)

$f_N \cdot f_+ \cdot P_+(0)$ and Λ_N
both decrease with B
and Λ_N is much too large
to be caused by either
 R_{HF} or neighbouring
nuclear dipoles.

Λ_C is also
anomalously fast.

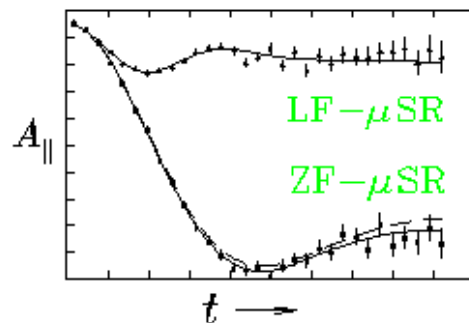
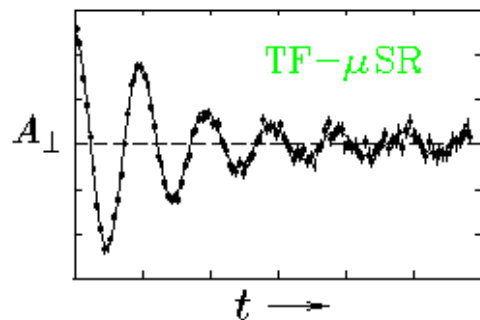


"Coulomb Explosion" Leftovers



Brewer's List of μ SR Acronyms

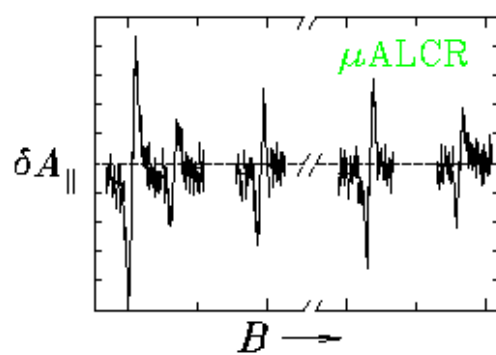
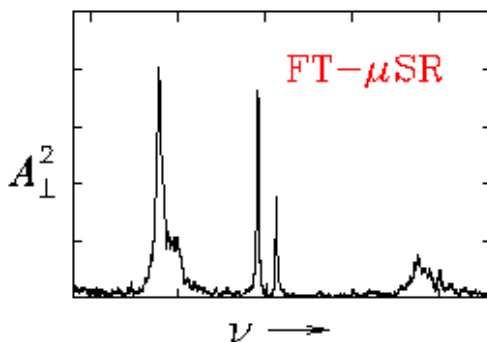
Transverse
Field



Longitudinal
Field

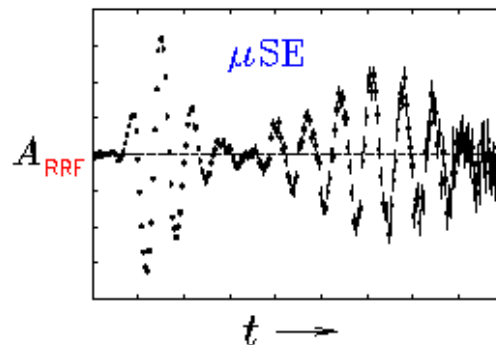
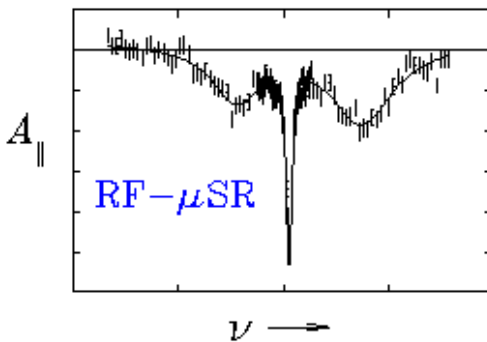
Zero Field

Fourier
Transform
 μ SR



Avoided
Level
Crossing
Resonance

Muon
Spin
Resonance



Muon
Spin
Echo